

Numerical solution of partial differential equations

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Course “Discontinuous Galerkin Method”

<https://www2.karlin.mff.cuni.cz/~dolejsi/Vyuka/DGM.html>

Partial differential equations

Why partial differential equations?

- many processes can be described (approximately) by PDEs
 - fluid dynamics, hydrology, heat and mass transfer, medicine, environmental protection, financial mathematics, etc.
 - these PDEs represent a mathematical description of physical, chemical, biological, etc. rules and/or laws
- some simplification usually necessary $\implies$ **model error**
- these PDEs are usually too complicated for an exact solution

Numerical solution of PDEs

- we solve PDEs approximately (numerically)
- we define new simplified (finite dimensional, solvable) problem $\implies$ **discretization error**

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Abstract problem described by PDEs

- let V be a functional space, we seek $u \in V$ such that
(EP) $\mathcal{L}u = f$
- $\mathcal{L}$ is a differential operator, f is a right-hand side,
- let solution of (EP) exists and is unique

Abstract numerical method

- let V_h be a space, $\dim(V_h) < \infty$, $V_h \subset V$ or $V_h \not\subset V$,
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Goals of the numerical solution of PDEs

Numerical analysis

- existence and uniqueness of u_h
- stability: if $\|u\| < \infty$ then $\|u_h\| < \infty$
- convergence: $u_h \rightarrow u$ for $\text{dof} = \dim(V_h) \rightarrow \infty$
- estimate $\|u - u_h\|$ in terms of dof (a priori estimate)
- estimate $\|u - u_h\|$ based on u_h (a posteriori estimate)
- robustness: validity of previous items for large range of data

Numerical realization

- algorithm for fast evaluation of u_h (efficiency)
- stability of the method in the finite precision arithmetic
- adaptive strategies = adaptive changes of V_h

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Numerical method

Numerical method in practise

- finite sequence of mathematical operations
- output is the approximate solution u_h

Construction of a numerical method for the given PDE

- discretization (space, time)
- setting of arising algebraic systems (numerical quadratures)
- (iterative) solution of nonlinear algebraic systems
- solution of linear algebraic systems

Type of discretizations

finite difference method, finite element method, finite volume method, spectral method, wavelets method, etc.

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- output is the approximate solution u_h

Construction of a numerical method for the given PDE

- **discretization** (space, time)
- setting of arising algebraic systems (numerical quadratures)
- (iterative) solution of nonlinear algebraic systems
- solution of linear algebraic systems

Type of discretizations

finite difference method, finite element method, finite volume method, spectral method, wavelets method, etc.

Choice of the numerical method

Which numerical method is the best one?

Depends on many aspects of the PDE considered

- physical background of the PDE
- expected regularity of the unknown exact solution
- presence of local phenomena
- outputs of interest
 - usual condition $\|u - u_h\| \leq TOL$ is not always practical
 - goal is the quantity of interest $J(u_h)$,
 $\Rightarrow$ error: $|J(u) - J(u_h)| \leq TOL$

Example

- conservation laws should be discretized by a conservative numerical method

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- **conservation laws** should be discretized by a **conservative numerical method**

Two basic physical processes

Diffusion

- $\frac{\partial u}{\partial t} - \nabla \cdot (a(u) \nabla u) = g$
- parabolic (elliptic) equation
- quantity is spread in all directions
- influence is decreasing for increasing distance of the source

Convection

- $\frac{\partial u}{\partial t} - \nabla \cdot \vec{f}(u) = g$
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- quantity is spread only in the direction of convection $\vec{f}(u)$
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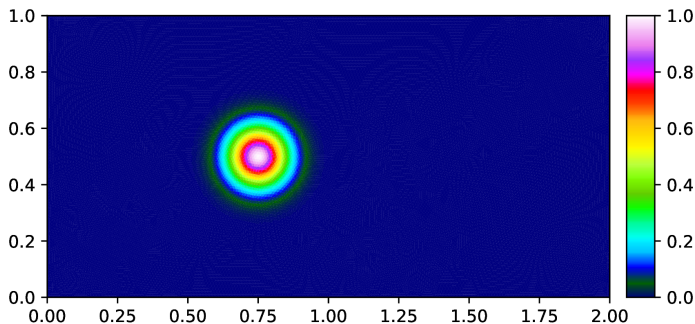
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Examples of physical features (1)

Only diffusion

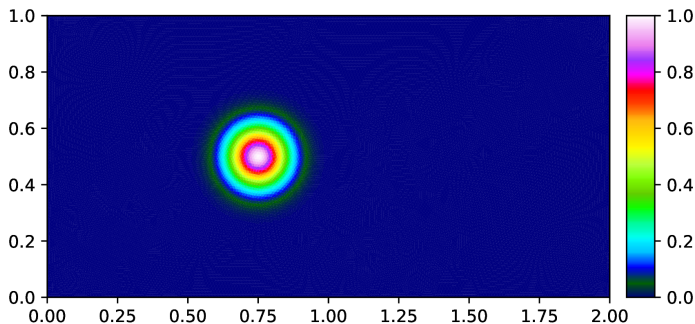
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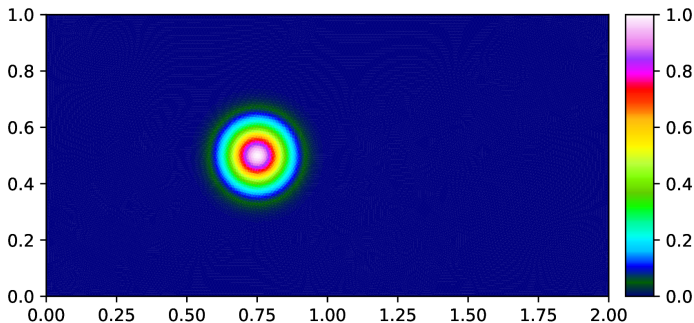
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Examples of physical features (2)

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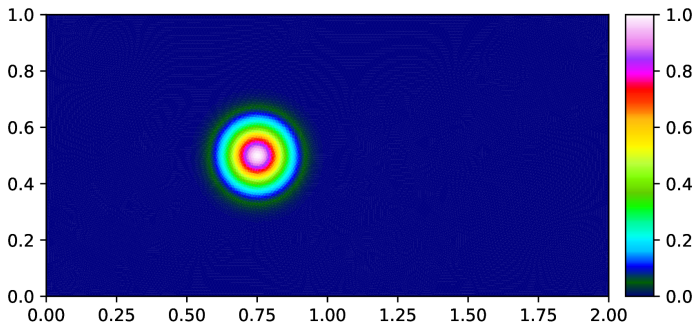
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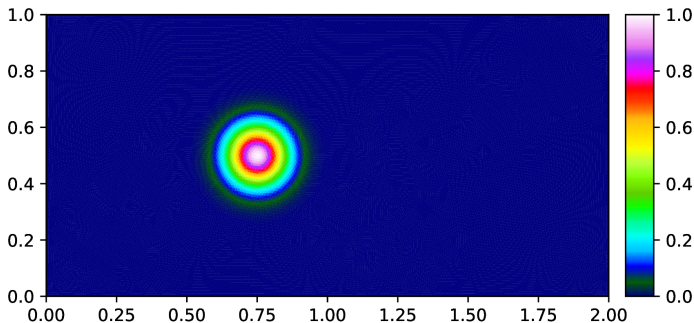
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Convection + small diffusion

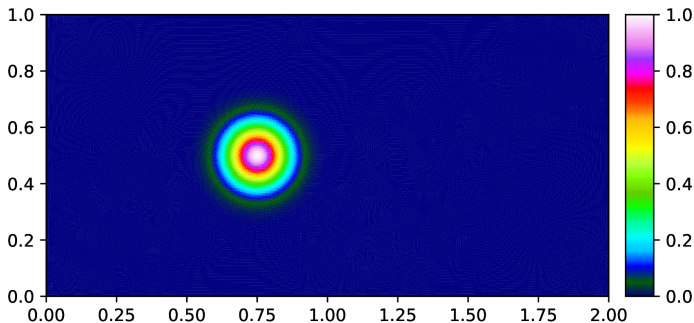
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Elliptic and parabolic PDE vs. hyperbolic PDE

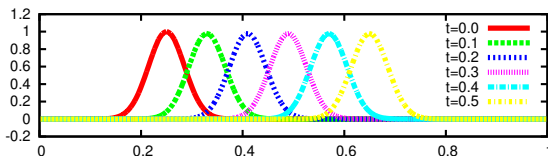
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$$u(x, 0) = \exp[(x - 1/4)^2]$$

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exact solutions for $\varepsilon = 0$, $\varepsilon = 10^{-4}$, $\varepsilon = 10^{-3}$, $\varepsilon = 10^{-2}$

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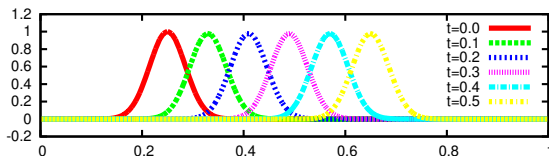
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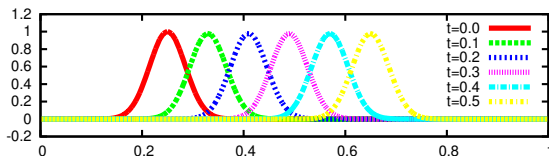
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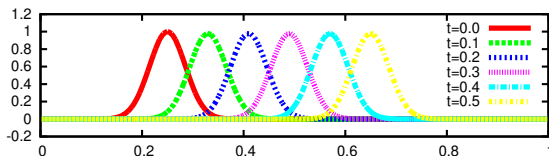
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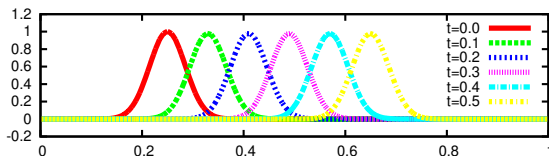
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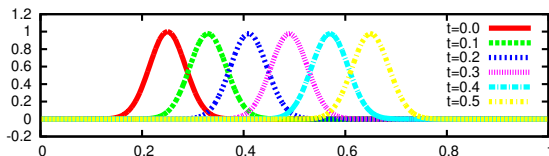
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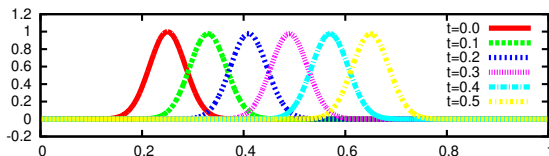
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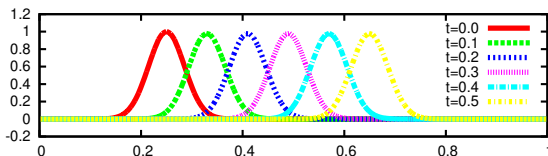
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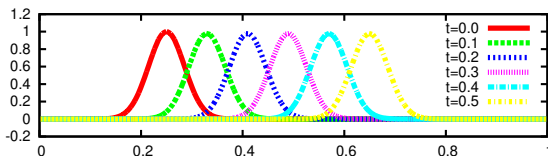
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Importance of the character of PDE

Why is important to know the previous properties?

- numerical solution is a kind of approximation
- many sources of inaccuracies:
 - discretization errors (finite dimensional approximation)
 - iterative errors (approximate solution of algebraic systems)
 - rounding errors (finite precision arithmetic)
- inaccuracies are propagated by PDEs and numerical scheme

Linear convection problem (no diffusion)

- exact solution: a simple propagation of the initial solution
- numerical solution: initial solution is propagated but smeared
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- numerical solution corresponds to convection+diffusion
- this effect is called numerical diffusion

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Why is important to know the previous properties?

- numerical solution is a kind of approximation
- many sources of inaccuracies:
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Effect of numerical diffusion

- zero diffusion does not exist in reality
- if numerical diffusion larger than physical one
 $\implies$ numerical solution can be completely wrong
- e.g., numerical solution is steady whereas reality is unsteady

Effect of “finite h ”

- we can prove that the proposed method is convergent
- approximate solution contains unphysical effects, e.g., spurious oscillations, negative temperature, etc.
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1D convection-diffusion equation

$$u : (0, 1) \rightarrow \mathbb{R} : \quad -\varepsilon u'' + u' = f, \quad u(0) = u(1) = 0, \quad \varepsilon > 0.$$

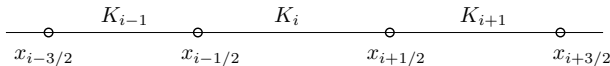
solution has a steep gradient near $x = 1$ (boundary layer)

Weak formulation

$$u \in H_0^1((0, 1)) : \quad \int_0^1 (\varepsilon u' v' + u' v) dx = \int_0^1 f v dx \quad \forall v \in H_0^1((0, 1))$$

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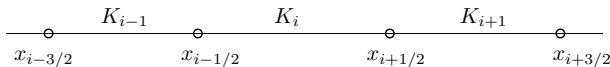
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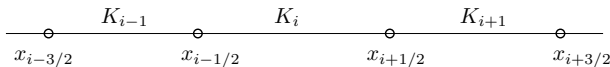
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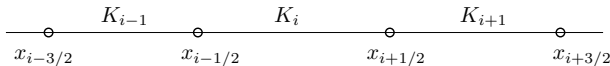
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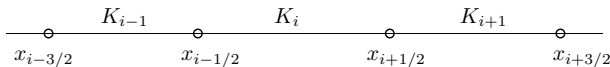
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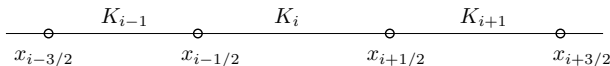
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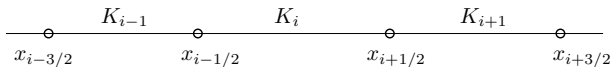
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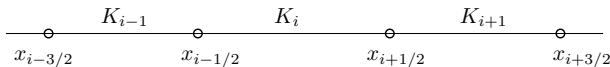
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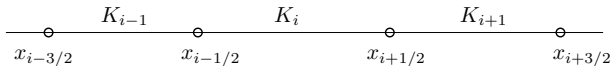
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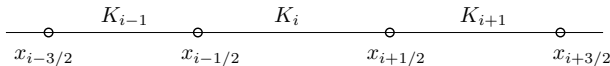
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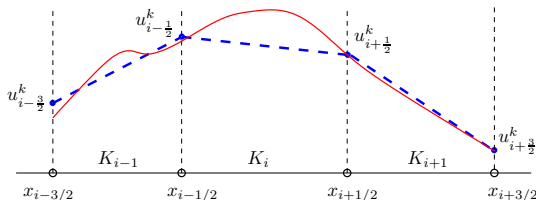
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Finite element method



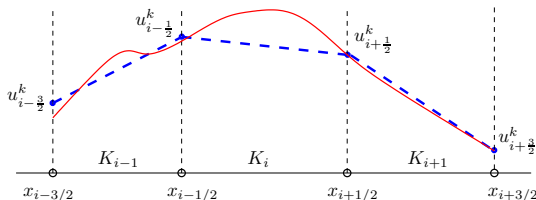
FEM solution

- $V_h = \{v_h \in C_0([0, 1]); v_h|_{K_i} = P^1(K_i), i = 1, \dots, N\}$
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$$\int_0^1 (\varepsilon u_h' v_h' + u_h' v_h) dx = \int_0^1 f v_h dx \quad \forall v_h \in V_h$$

- reasonable discretization of diffusion $\Rightarrow$ we prove convergence
- discretization of convective term “does not respect physics”

Finite element method



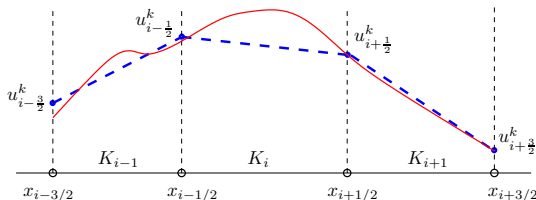
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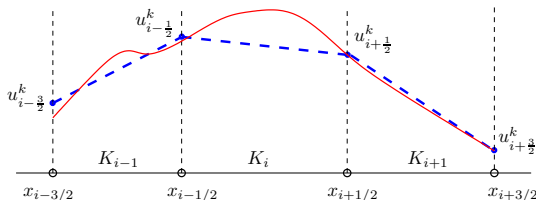
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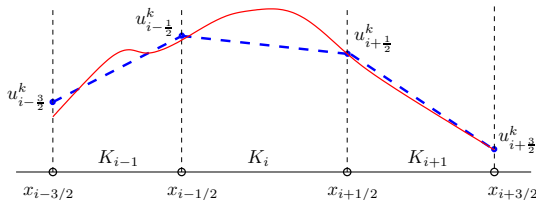
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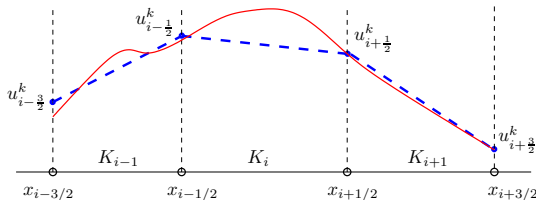
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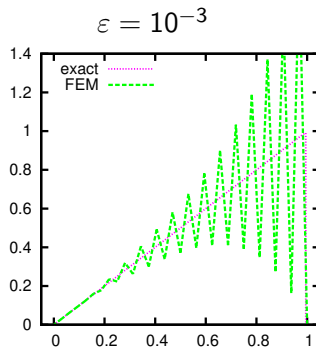
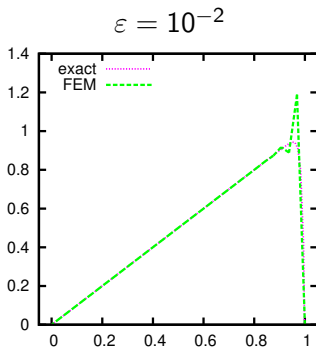
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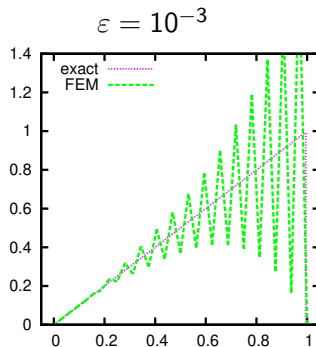
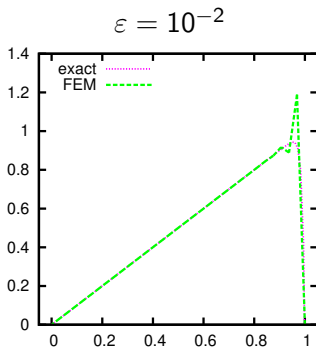
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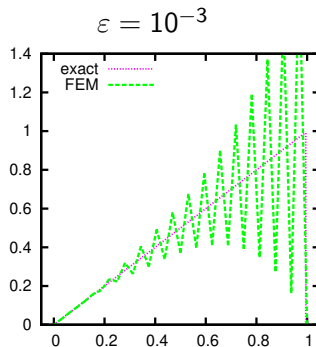
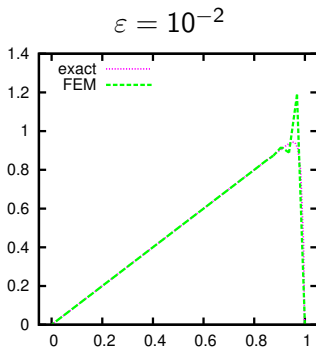
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Finite element method



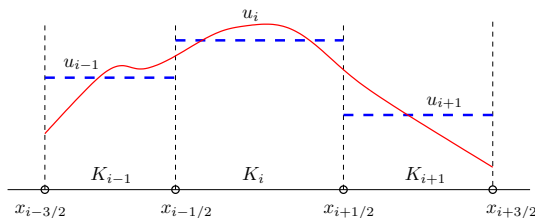
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Finite volume method



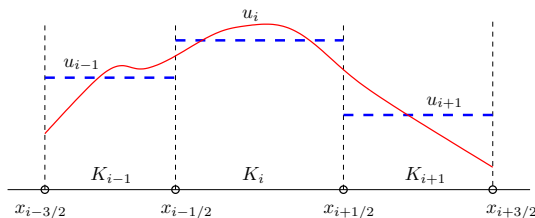
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- $V_h = \{v_h \in L^2([0, 1]); v_h|_{K_i} = P^0(K_i), i = 1, \dots, N\}$
- we integrate $-\varepsilon u'' + a u' = 1$ over K_i and use Gauss theorem

$$-\varepsilon[u'(\cdot)]_{x_{i-1/2}}^{x_{i+1/2}} + a[u(\cdot)]_{x_{i-1/2}}^{x_{i+1/2}} = |K_i|$$

- $u|_{x_{i+1/2}} = ??$ upwinding: $a > 0 \Rightarrow u|_{x_{i+1/2}} := u_i$

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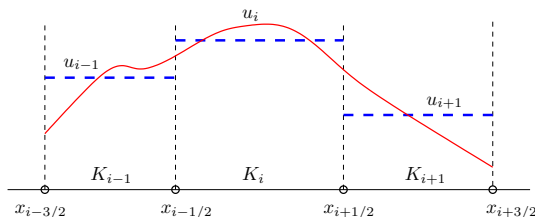
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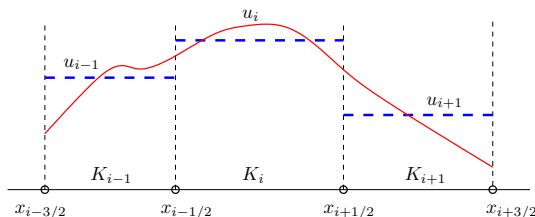
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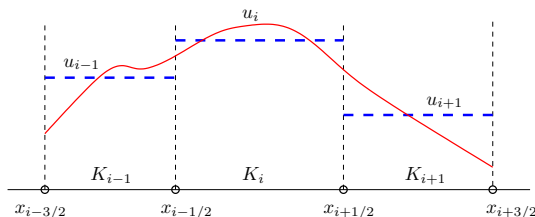
Piecewise constant approximation

- $V_h = \{v_h \in L^2([0, 1]); v_h|_{K_i} = P^0(K_i), i = 1, \dots, N\}$
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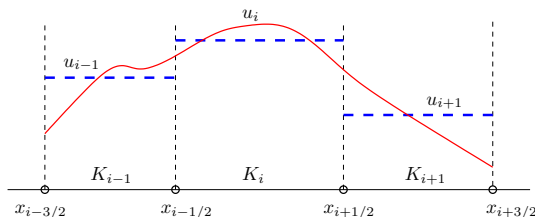
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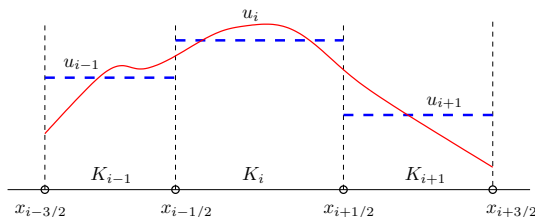
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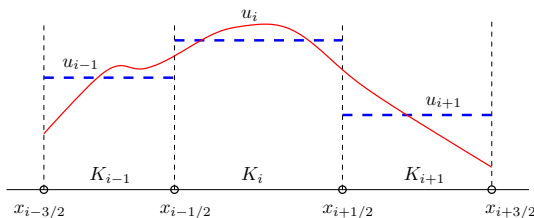
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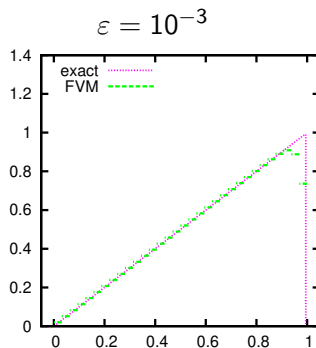
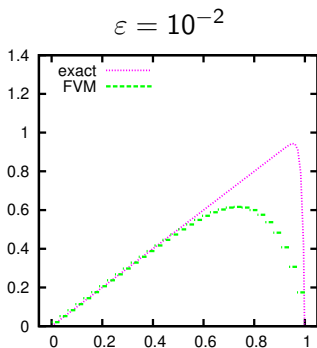
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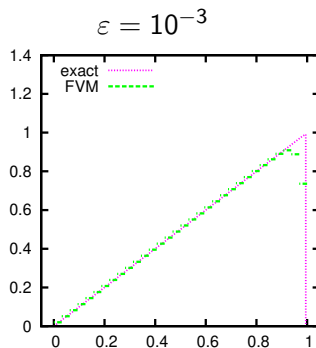
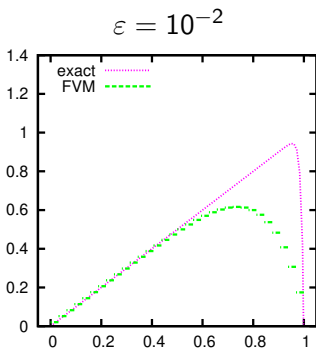
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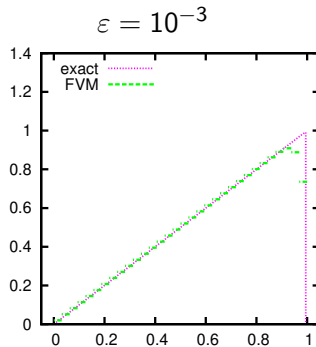
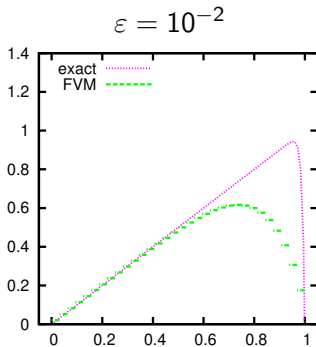
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- Low accuracy for larger ε
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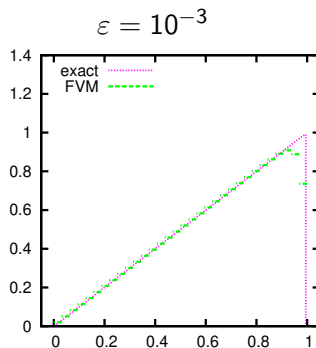
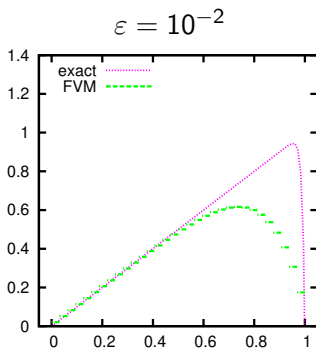
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comparison of FEM and FVM for **time-dependent convective** problem

Finite element method

- continuous approximation
- high order of accuracy
- many theoretical results
- fine for diffusive problems

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Discontinuous Galerkin method

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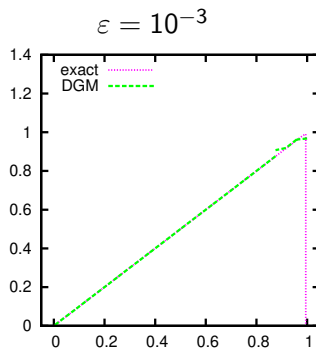
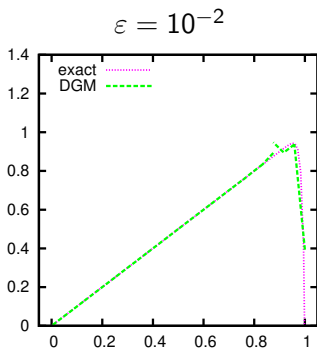
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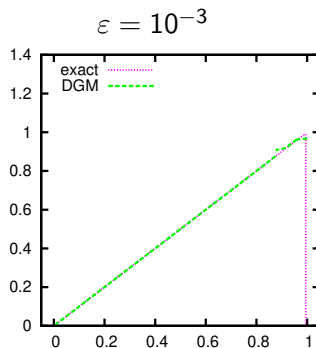
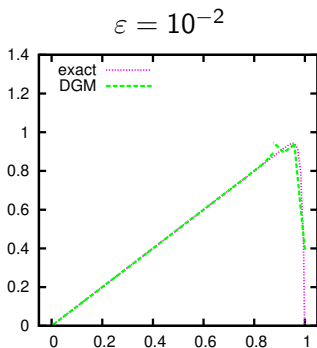
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- not ideal but works very well for both ε
- additional techniques (remedies) are possible

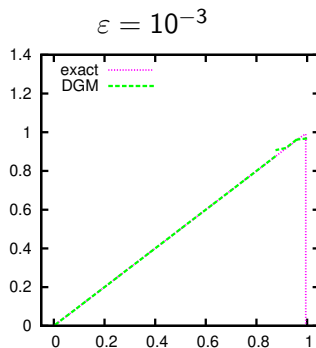
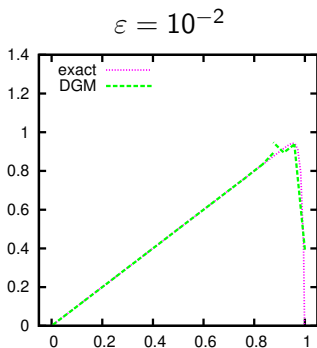
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- efficient method for the numerical solution of various PDEs
- piecewise polynomial BUT discontinuous approximation
- suitable for very large range of problems
 - elliptic, parabolic, hyperbolic
 - linear, nonlinear, degenerate
- space-time DGMs are available
- flexibility in the mesh design
 - non-matching and non-uniform grids
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 - varying polynomial approximation degrees
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- numerical analysis of the method is more complicated

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- Abstract error analysis
- DGM for the Laplace problem: complete error analysis
- numerical approximation based on upwinding
- DGM for the nonlinear convection-diffusion equation
- DGM for time dependent problems
- DGM for compressible flow problems and other applications

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- standard lectures (lecture notes are available)
- 3 quizzes during the semester (can replace the exam)
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- DGM for time dependent problems
- DGM for compressible flow problems and other applications

Organization issues

- standard lectures (lecture notes are available)
- 3 quizzes during the semester (can replace the exam)
- oral exam

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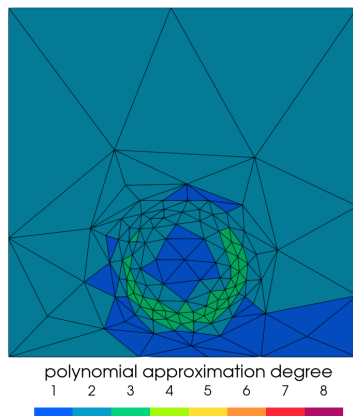
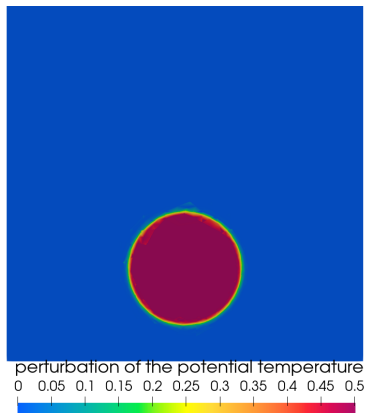
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Modelling of mesoscale atmosphere by DGM

compressible Navier-Stokes equations with the gravity term



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