

Department of Mechanical Engineering
Indian Institute of Technology Gandhinagar

Dynamic Fragmentation of Functionally Graded Brittle Materials in 1D

Ravi Sastri Ayyagari

[Work of my PhD Student: Archita Gogoi]

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Introduction

Methodology

- 1D FGMs under uniform strain rates: Displacement fields
- Decomposition of Displacement Field
- Boundary Conditions
- Minimum fragment size

Results

› Fragmentation

A contiguous body breaks into several pieces when loaded beyond design capacity.

› Application

Fragments are desired (e.g., kidney stones), prevented (e.g., bulletproof vests), or controlled (e.g., rock mining).

› Analytical Modelling

Predicts minimum fragment size in materials under constant dynamic strain rates.

Q Can fragmentation in brittle materials be tailored?

Functionally Graded Materials

Spatial variation in the material properties.

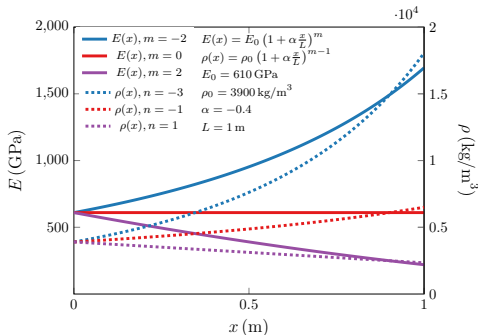
Engineering Applications

High fidelity & robustness.

Testing & Characterization

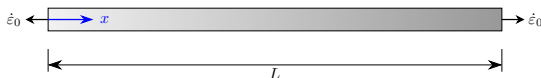
Elastic wave propagation in FGM's

- » Reliable NDT methods
- » Assess structural integrity



Q Can we design FGMs with tailored wave propagation characteristics to control fragmentation behaviour?





› Governing Equation

$$\left(\hat{E}(x) u_{,x} \right)_{,x} = \rho(x) u_{,tt} \quad \text{where} \quad c(x) = \sqrt{\hat{E}(x)/\rho(x)}$$

› Initial Conditions

$$u(x, 0) = 0, \quad u_{,t}(x, 0) = \dot{\epsilon}_0 \left(x - \frac{L}{2} \right)$$

› Displacement Field (Transcendental Form)

For material properties exhibiting a polynomial variation:

$$\hat{U}(z, s) = C_1 I_0 \left(\frac{2s}{\alpha} \sqrt{z} \right) + C_2 K_0 \left(\frac{2s}{\alpha} \sqrt{z} \right) + \frac{\dot{\epsilon}_0}{s^2 \alpha} z + \frac{\dot{\epsilon}_0}{s^4 \alpha} (\alpha^2 - s^2)$$

where $z = \left(1 + \frac{\alpha x}{L} \right)$, $I_0 \left(\frac{2s}{\alpha} \sqrt{z} \right)$, and $K_0 \left(\frac{2s}{\alpha} \sqrt{z} \right)$ are the modified Bessel functions of the first and second kinds, respectively.

Q Can we get a closed-form displacement field $u(x, t)$?

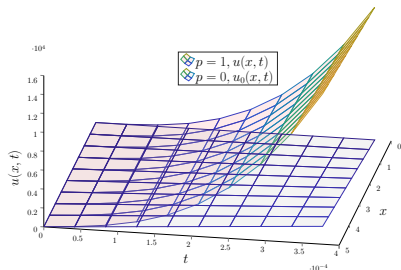
Let $u = u_0 + pu_1 + p^2u_2 + \dots$

➤ Governing Equation:

$$u_{,tt} = p\{c_0^2(1 + \frac{\alpha x}{L})u_{,xx} + c_0^2 \frac{\alpha x}{L}u_{,x}\}$$

where:

$$u = \lim_{p \rightarrow 1} (u_0 + pu_1 + p^2u_2 + \dots)$$



Form	$\hat{E}(x) = E_0 \times$	$\rho(x) = \rho_0 \times$	Cases	$u(x, t)$
Polynomial	ψ^m	ψ^n	$m - n = 1$	$\dot{\epsilon}_0 t (x - L/2 + m\phi)$
			$m - n = 2$	$\frac{\dot{\epsilon}_0 \psi L^2}{\alpha^2 \sqrt{m} c_0} \sinh\left(\frac{\sqrt{m} c_0 \alpha t}{L}\right) - \frac{\dot{\epsilon}_0 t L}{\alpha}$
Sinusoidal	$\sin^m(\beta(x) - \alpha)$	$\sin^n(\beta(x) - \alpha)$	$\begin{matrix} m - n = 0 \\ (m = n = 2) \end{matrix}$	$\dot{\epsilon}_0 t (x - L/2 + m\phi \cot(\beta(x) - \alpha))$
Exponential	$\exp(m\beta(x))$	$\exp(n\beta(x))$	$m - n = 0$	$\dot{\epsilon}_0 t (x - L/2 + m\phi)$

† Above $\beta(x) = x/L$. (E_0, ρ_0, c_0) correspond to properties at $x = 0$.

‡ α is an arbitrary real constant, $\psi = (1 + \alpha\beta(x))$ and $\phi = c_0^2 \alpha t^2 / (6L)$.

Q What are the forward and backward components?

Objective : To determine the forward and backward components of the displacement field.

› **Step 1:** Let $u_t = \hat{v}$ and $u_x = \hat{w}$ to get

$$U_t + \mathcal{D}(x, t)U_x = \mathbf{F}(x, t)$$

where: $U = \begin{Bmatrix} u_t \\ u_x \end{Bmatrix}$, $\mathcal{D} = \begin{bmatrix} 0 & -c^2(x) \\ -1 & 0 \end{bmatrix}$, and $\mathbf{F} = \begin{Bmatrix} \frac{\hat{E}_x(x)}{\rho(x)}\hat{w} \\ 0 \end{Bmatrix}$

› **Step 2:** Use the linear transformation $U = \mathcal{P}(x, t)V$ such that

$$u_t = v_1 + v_2, \quad v_{1,2} = \left(\frac{u_t \pm c(x)u_x}{2} \right) \quad \text{where: } V = \begin{Bmatrix} v_1 \\ v_2 \end{Bmatrix}$$

› **Step 3:** From $U = \mathcal{P}(x, t)V$, we diagonalize $\mathcal{D}(x, t)$ to get

$$V_t + \begin{bmatrix} -c(x) & 0 \\ 0 & c(x) \end{bmatrix} V_x = \mathbf{H}(x, t)$$

Objective : To determine the forward and backward components of the displacement field.

› **Step 1:** Let $v_t = \hat{E}(x)u_x$ and $v_x = \rho(x)u_t$ to get

$$U_t + \mathcal{D}(x, t)U_x = \mathbf{F}(x, t) \quad \text{Non-unique !}$$

where: $U = \begin{Bmatrix} u \\ v \end{Bmatrix}$, $\mathcal{D} = \begin{bmatrix} 0 & -\frac{1}{\rho(x)} \\ -\hat{E}(x) & 0 \end{bmatrix}$, and $\mathbf{F} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$

› **Step 2:** Use the linear transformation $U = \mathcal{P}(x, t)V$ such that

$$u = v_1 + v_2, \quad v_{1,2} = \left(\frac{uZ(x) \pm v}{2Z(x)} \right) \quad \text{where: } V = \begin{Bmatrix} v_1 \\ v_2 \end{Bmatrix}$$

and $Z(x) = \sqrt{\hat{E}(x)\rho(x)}$

› **Step 3:** From $U = \mathcal{P}(x, t)V$, we diagonalize $\mathcal{D}(x, t)$ to get

$$V_t + \begin{bmatrix} -c(x) & 0 \\ 0 & c(x) \end{bmatrix} V_x = \mathbf{H}(x, t)$$

when $\mathbf{H}(x, t) \neq 0$, v_1 and v_2 are **Riemann Variables**.

Objective : To determine the forward and backward components of the displacement field.

‣ **Step 1:** Let $v_t = \hat{E}(x)u_x$ and $v_x = \rho(x)u_t$ to get

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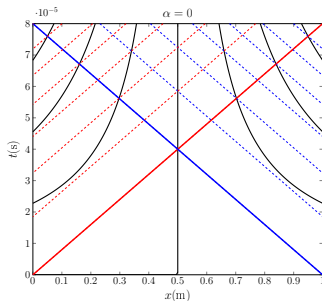
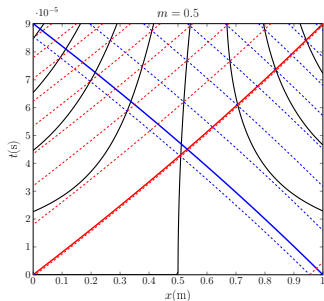
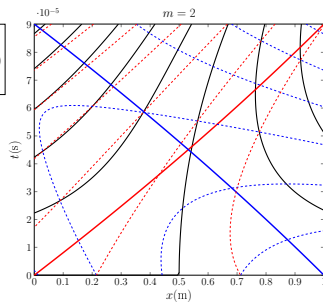
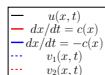
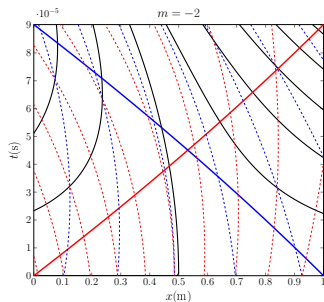
$$V_t + \begin{bmatrix} -c(x) & 0 \\ 0 & c(x) \end{bmatrix} V_x = \mathbf{H}(x, t)$$

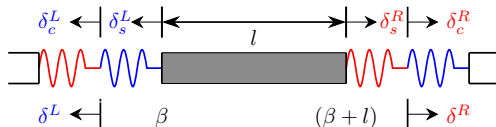
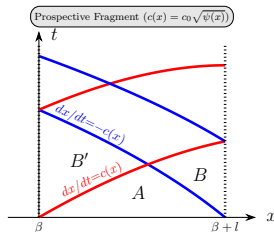
when $\mathbf{H}(x, t) = 0$, v_1 and v_2 are **Riemann Invariants**.

Decomposition of Displacement Field - II



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$$u(\beta, t) = \dot{\epsilon}_0(\beta - L/2)t + \delta_c^L(t)$$

$$E(\beta) u_x(\beta, t) = \sigma_c(\delta_c(t))$$

$$u(\beta + l, t) = \dot{\epsilon}_0(\beta + l - L/2)t - \delta_c^R(t)$$

$$E(\beta + l) u_x(\beta + l, t) = \sigma_c(\delta_c(t))$$

where $\delta_c^L(t) + \delta_c^R(t) = \delta_c(t)$ and $\sigma_c(\delta_c(t))$ is a suitable cohesive law.

- From the boundary conditions and after non-dimensionalising:

$$\tilde{\delta}_c(\tilde{t}) + \tilde{\delta}_c(\tilde{t}) \exp(1 - \tilde{\delta}_c(\tilde{t})) = \tilde{\epsilon}_0 \tilde{t} + \tilde{\epsilon}_0 \beta_1$$

where

$$\beta_1 = \frac{\tilde{l}(-3 + 2(m + 1))}{4(1 + m)} \left(1 + \frac{\epsilon^2 + C}{\epsilon(C + 1)} \right), \quad \epsilon = \frac{E(\tilde{\beta} + \tilde{l})}{E(\tilde{\beta})}, \quad C = \frac{c(\tilde{\beta} + \tilde{l})}{c(\tilde{\beta})}$$

- We numerically solve this ODE to obtain values of $\tilde{\delta}_c(\tilde{t})$.
- Conditions for fragmentation to occur:

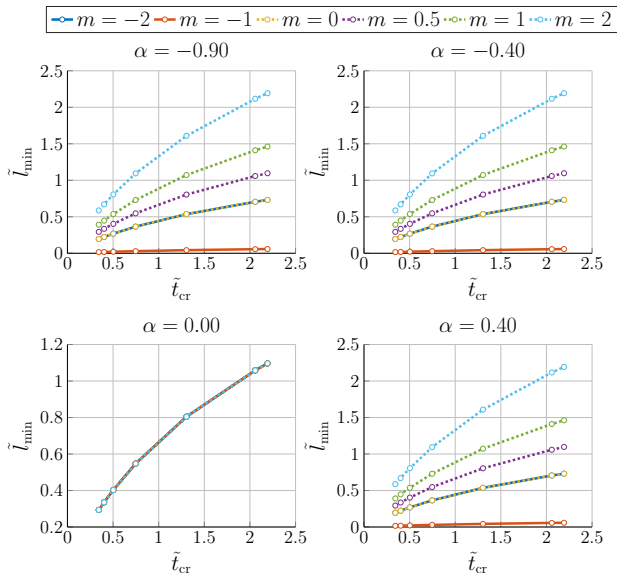
$$\delta_c(t_{\text{cr}}) = \delta^* \implies \tilde{\delta}_c(\tilde{t}_{\text{cr}}) = 1 \quad (1)$$

$$\dot{\delta}_s(t_{\text{cr}}) = 0 \implies \dot{\delta}_c(t_{\text{cr}}) = \dot{\epsilon}_0 l_{\min} \implies \tilde{\delta}_c(\tilde{t}_{\text{cr}}) = \tilde{\epsilon}_0 \tilde{l}_{\min} \quad (2)$$

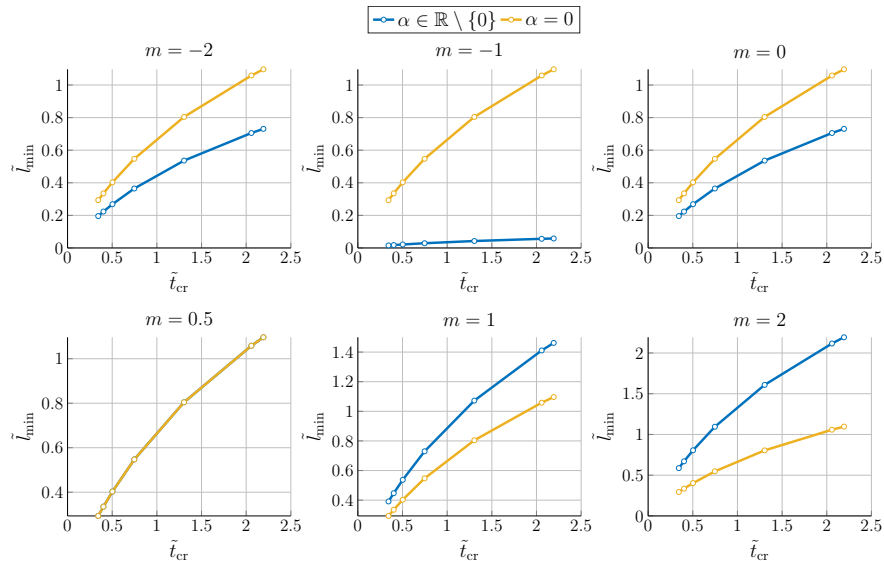
- The conditions are substituted in the ODE to obtain an implicit equation in terms of $\tilde{l}_{\min}$.
- This implicit equation is solved to obtain

$$\tilde{l}_{\min} = f(\alpha, m, \beta, L, \dot{\epsilon}_0, \tilde{t}_{\text{cr}})$$

Dependence on m :



Dependence on α :



- › Fragment size distribution
- › Number of fragments
- › Robust methods for decomposition of displacement fields even when the field is described in terms of special functions
- › Fragmentation in higher dimensions
- › Fragmentation under multi-axial loading

Thank you !

Questions ?