

Modelling, simulation and benchmarking of fluid-structure interactions with contact



Stefan Frei

Modelling, PDE analysis and computational mathematics in materials science
Prag, 25.09.2024

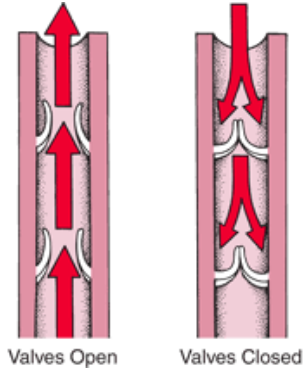
Joint work with

E. Burman (UCL London), M.A. Fernández (Inria Paris), T. Richter (OvGU Magdeburg)

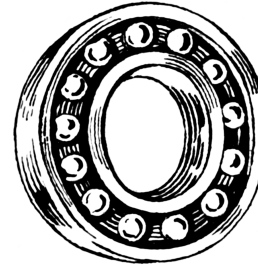
Motivation

Typically **fluid** or **gas** between **contacting structures** before/after contact

Vessel valves

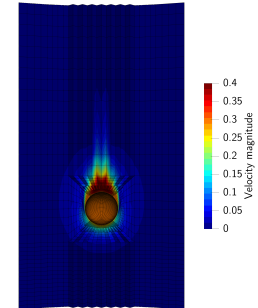
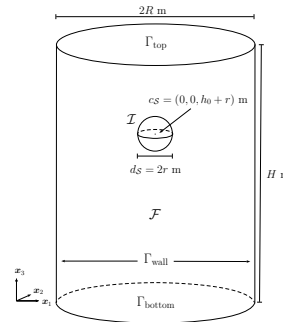


Ball bearing (Lubricant)



Experimental benchmark: Falling and bouncing **elastic balls** (PTFE, rubber) in a **water-glycerine** mixture

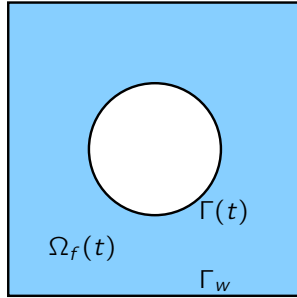
(HAGEMEIER, THEVENIN, RICHTER, Int J Multiphase Flow (2020)
VON WAHL, RICHTER, F., HAGEMEIER, Phys Fluids (2021))



Overview

- 1 Fluid-structure interaction and contact
- 2 A model for seepage
- 3 Discretisation and numerical results

Fluid-structure interaction



Fluid problem

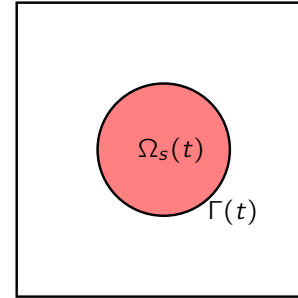
Incompr. Navier-Stokes equations in $\Omega_f(t)$

$$\begin{aligned}\rho_f(\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u}) - \operatorname{div} \sigma_f(\mathbf{u}, p) &= \mathbf{f}_f, \\ \operatorname{div} \mathbf{u} &= 0\end{aligned}$$

Boundary conditions (no-slip)

$$\mathbf{u} = 0 \quad \text{on } \Gamma_w \quad + \text{ other bc}$$

No-collision paradox: The Navier-Stokes equations with **no-slip** boundary or interface conditions do not allow for contact (HILLAIRET/HESLA 2D; GERART-VARET, HILLAIRET AND WANG 3D)



Solid problem

Linear elastic material in $\Omega_s(t)$

$$\begin{aligned}\rho_s(\partial_t \mathbf{d} + \mathbf{d} \cdot \nabla \mathbf{d}) - \operatorname{div} \sigma_s(\mathbf{d}) &= \mathbf{f}_s, \\ \partial_t \mathbf{d} + \mathbf{d} \cdot \nabla \mathbf{d} &= \dot{\mathbf{d}}\end{aligned}$$

Coupling conditions (no-slip)

$$\mathbf{u} = \dot{\mathbf{d}}, \quad \sigma_f \mathbf{n} = \sigma_s \mathbf{n} \quad \text{on } \Gamma(t).$$

Relaxed contact formulation

We impose the contact condition w.r.t to Γ_ϵ , such that a **small fluid layer** of size $\epsilon(h)$ remains

- The **"no-penetration"** condition reads

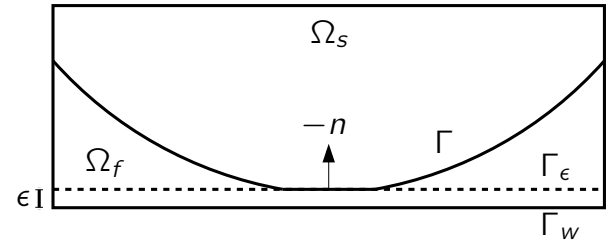
$$d_n \leq g_\epsilon := g_0 - \epsilon(h).$$

- The contact conditions are formulated with a Lagrange multiplier λ

$$d_n \leq g_\epsilon, \quad \lambda \leq 0, \quad (d_n - g_\epsilon)\lambda = 0$$

- Elimination of the Lagrange multiplier ($\lambda = \underbrace{\sigma_{s,nn} - \sigma_{f,nn}}_{[\sigma_{nn}]}$)

$$d_n \leq g_\epsilon, \quad [\sigma_{nn}] \leq 0, \quad (d_n - g_\epsilon)[\sigma_{nn}] = 0 \quad (1)$$



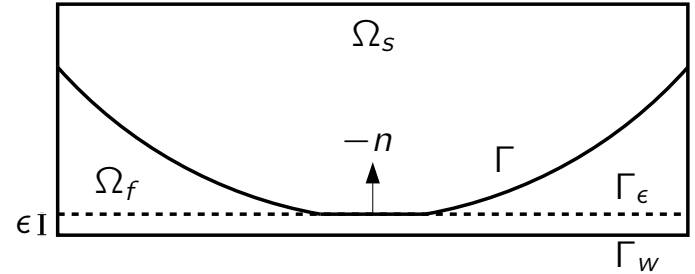
For arbitrary $\gamma > 0$ (1) is equivalent to (ALART & CURNIER '91)

$$[\sigma_{nn}] = -\gamma \max \left\{ d_n - g_\epsilon - \frac{1}{\gamma} [\sigma_{nn}], 0 \right\} =: -\gamma [P_\gamma([\sigma_{nn}], d_n)]_+ \quad \text{on } \Gamma(t).$$

Problems

Relaxed contact formulation

$$d_n \leq g_\epsilon, \quad \llbracket \sigma_{nn} \rrbracket \leq 0, \quad (d_n - g_\epsilon) \llbracket \sigma_{nn} \rrbracket = 0$$



- The **fluid forces** $\sigma_{f,nn}$ in the **fluid layer** enter the contact dynamics
- As in particular the **pressure** p might take large values, these can have a **non-physical impact** on the contact dynamics
- We need a **physically motivated model** in the layer!

Idea:

- The layer appears due to **surface roughness** on a **micro scale** of thickness ϵ_p . We formulate a **Darcy law** in the layer and take $\epsilon_p \rightarrow 0$
- Similar ideas with a porous domain of thickness $\epsilon_p > 0$: Complex Biot-type equations (AGER, SCHOTT, VUONG, POPP & WALL, 2019), **Navier-Stokes-Brinkmann** (GEROSA & MARSDEN, 2024)

Navier-Stokes-Darcy coupling

- In the porous domain Ω_p , we assume a **Darcy law**

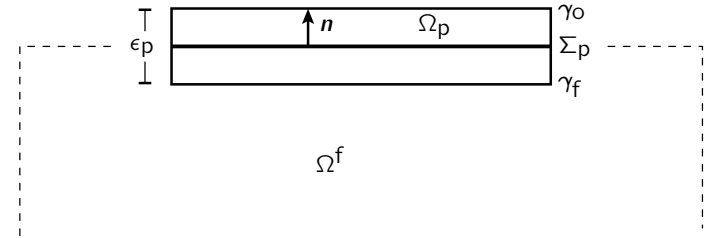
$$\begin{cases} u_l + \mathbf{K} \nabla p_l = 0 & \text{in } \Omega_p, \\ \nabla \cdot u_l = 0 & \text{in } \Omega_p, \end{cases}$$

where $\mathbf{K} \in \mathbb{R}^{d \times d}$ with

$$\mathbf{K} \nabla p_l = K_\tau \nabla_\tau p_l + K_n \partial_n p_l.$$

- These equations are coupled to **Navier-Stokes** by means of the **Beavers-Joseph-Saffman** conditions

$$\begin{cases} \sigma_{f,n} = -p_l & \text{on } \gamma_f, \\ u_n = u_l \cdot \mathbf{n} & \text{on } \gamma_f, \\ \sigma_{f,n\tau} = -\frac{\alpha}{\sqrt{K_\tau \epsilon_p}} u_\tau & \text{on } \gamma_f \end{cases}$$



- For $\epsilon_p \rightarrow 0$, we obtain on the mid-surface Σ_p (MARTIN ET. AL, 2005)

$$-\nabla_\tau \cdot (\epsilon_p K_\tau \nabla_\tau P_l) = u_n \quad \text{on } \Sigma_p,$$

$$\sigma_{f,nn} = -P_l - \frac{\epsilon_p K_n^{-1}}{4} u_n \quad \text{on } \Sigma_p,$$

$$\sigma_{f,n\tau} = -\frac{\alpha}{\sqrt{K_\tau \epsilon_p}} u_\tau \quad \text{on } \Sigma_p$$

where $P_l := \frac{1}{2} (p_l|_{\gamma_f} + p_l|_{\gamma_0})$.

Complete model

- **Solid problem** in $\Omega_s(t)$:

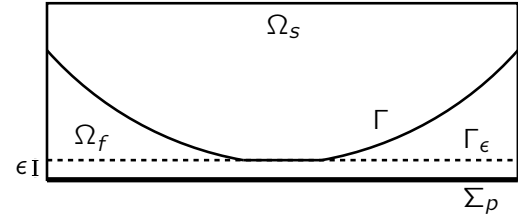
$$\begin{cases} \rho_s(\partial_t \dot{d} + \dot{d} \cdot \nabla \dot{d}) - \operatorname{div} \boldsymbol{\sigma}_s(d) = \rho_s \mathbf{f}_s, \\ \partial_t d + \dot{d} \cdot \nabla d = \dot{d}, \\ d = 0 \text{ on } \Gamma_s^d, \end{cases}$$

- **Fluid problem** in $\Omega_f(t)$

$$\begin{cases} \rho_f(\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u}) - \operatorname{div} \boldsymbol{\sigma}_f(\mathbf{u}, p) = \rho_f \mathbf{f}_f, \\ \operatorname{div} \mathbf{u} = 0, \\ \mathbf{u} = 0 \text{ on } \Gamma_f^d \end{cases}$$

- **Porous layer** (seepage)

$$\begin{cases} -\nabla_\tau \cdot (\epsilon_p K_\tau \nabla_\tau P_l) = u_n & \text{on } \Sigma_p, \\ \epsilon_p K_\tau \tau \cdot \nabla_\tau P_l = 0 & \text{on } \partial \Sigma_p, \end{cases}$$



- **Fluid-porous coupling conditions:**

$$\begin{cases} \sigma_{f,nn} = -P_l - \underbrace{\frac{\epsilon_p K_n^{-1}}{4} u_n}_{=: \sigma_p} & \text{on } \Sigma_p, \\ \sigma_{f,n\tau} = -\frac{\alpha}{\sqrt{K_\tau \epsilon_p}} u_\tau & \text{on } \Sigma_p. \end{cases}$$

- We combine **FSI** and **contact conditions** on the joint interface-contact surface $\Gamma(t)$ to

$$u = \dot{d}, \quad \llbracket \sigma_{nn} \rrbracket = -\gamma [P_\gamma(\llbracket \sigma_{nn} \rrbracket, d_n)]_+$$

Variational formulation (Fully Eulerian)

Simultaneous imposition of FSI interface and contact conditions on a **joint interface-contact surface** $\Gamma(t)$ by means of Nitsche's method (BURMAN, FERNÁNDEZ, F., GEROSA, CMAME 2022)

Find $u(t) \in \mathcal{V}_h$, $p(t) \in \mathcal{Q}_h$, $d(t)$, $\dot{d} \in \mathcal{W}_h$, $P_l \in \mathcal{S}_h$, such that

$$\begin{aligned}
 & a_f(u, p; v, q) + a_s(d, \dot{d}; w, z) - (\sigma_f(u, p)\mathbf{n}, w - v)_{\Gamma(t)} - (\dot{d} - u, \sigma_f(v, -q)\mathbf{n})_{\Gamma(t)} \\
 & + \frac{\gamma_{\text{fsi}}}{h} (\dot{d} - u, w - v)_{\Gamma(t)} + \frac{\gamma_c}{h} ([P_\gamma([\sigma_{nn}], d_n)]_+, w_n)_{\Gamma(t)} \\
 & - (\sigma_p, v_n)_{\Sigma_p} + (\epsilon_p K_\tau \nabla_\tau P_l, \nabla_\tau q_l)_{\Sigma_p} - (u_n, q_l)_{\Sigma_p} + \frac{\alpha}{\sqrt{K_\tau \epsilon_p}} (u_\tau, v_\tau)_{\Sigma_p} \\
 & = (f_f, v)_{\Omega_f(t)} + (f_s, w)_{\Omega_s(t)} \quad \forall v, q, w, z, q_l \in \mathcal{V}_h \times \mathcal{Q}_h \times \mathcal{W}_h \times \mathcal{W}_h \times \mathcal{S}_h
 \end{aligned}$$

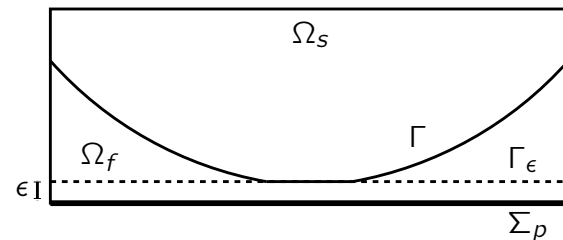
where

$$\begin{aligned}
 a_f(u, p; v, q) & := \rho_f (\partial_t u + u \cdot \nabla u, v)_{\Omega_f(t)} + (\sigma_f(u, p), \nabla v)_{\Omega_f(t)} + (\text{div } u, q)_{\Omega_f(t)} \\
 a_s(d, \dot{d}; w, z) & := (\rho_s (\partial_t \dot{d} + \dot{d} \cdot \nabla \dot{d}), w)_{\Omega_s(t)} + (\sigma_s(d), \nabla w)_{\Omega_s(t)} + (\partial_t d + \dot{d} \cdot \nabla d - \dot{d}, z)_{\Omega_s(t)}
 \end{aligned}$$

Advantages & mechanical consistency

The **relaxed contact formulation** greatly simplifies the implementation:

- The coupling is for all times **structure-fluid** and **fluid-porous** medium (never **structure-porous** medium)
- Topology changes are avoided



Collisions are possible, even when no-slip is used on the FSI interface

(CHAMPION, FERNÁNDEZ, GRANDMONT, VERGNET, VIDRASCU 2024)

Mechanical consistency:

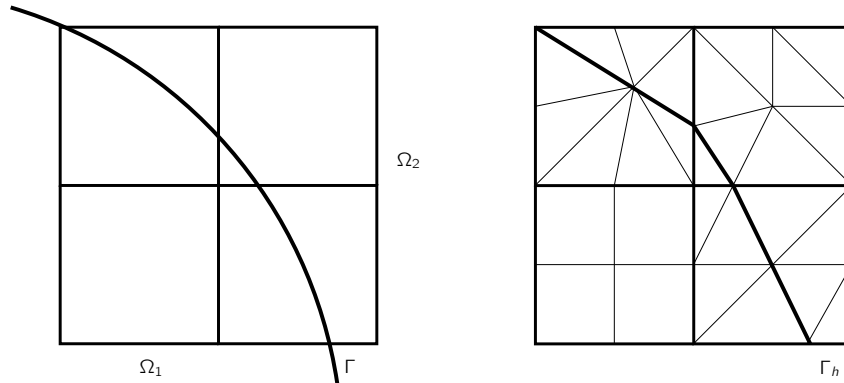
- If a part of $\Gamma(t)$ is in '**contact**' with Σ_p (in the relaxed sense), it holds

$$\sigma_{f,nn}|_{\Gamma} \approx \sigma_p|_{\Sigma_p}.$$

- This gives a **physical meaning** to the **fluid forces** $\sigma_{f,nn}$ in the layer.

Discretisation: Fitted (locally modified) finite elements

- Fixed regular **patch mesh** independent of the interface location
- Split interface patches into eight triangles to **resolve the interface**
- Combination of P_1 and Q_1 finite elements (F. & RICHTER, SINUM 2014)



- Equal-order locally mod. FE with anisotropic **edge-oriented pressure stabilisation** (F., IJNMF, 2019)
- Time discretisation: **Modified dG(0)** scheme (F., RICHTER, M2AN, 2017)

Numerical example

- Fall of an elastic **PTFE ball** within a **water-glycerin mixture** (2d)

- **Porous medium**

$$\epsilon_p = 10^{-4}, K_T = K_n = 10^{-2}$$

- **Contact** parameters

$$\gamma_c = 30\lambda_s, \epsilon_g = \frac{h_{\min}}{4}$$

- **Adaptive time-stepping**

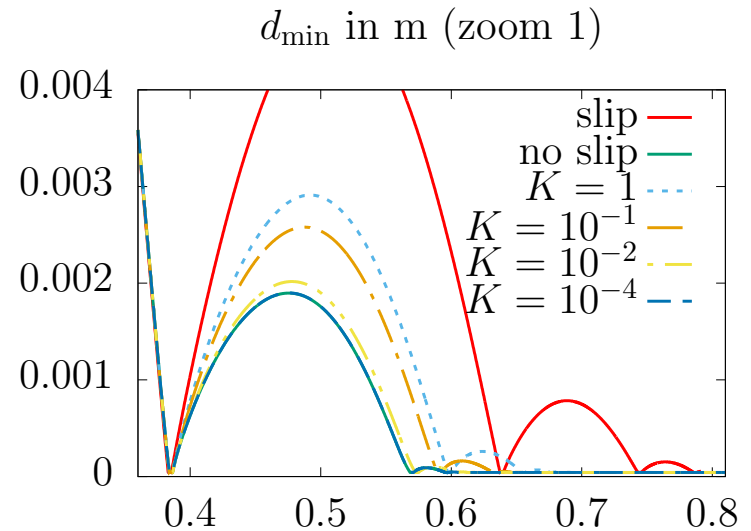
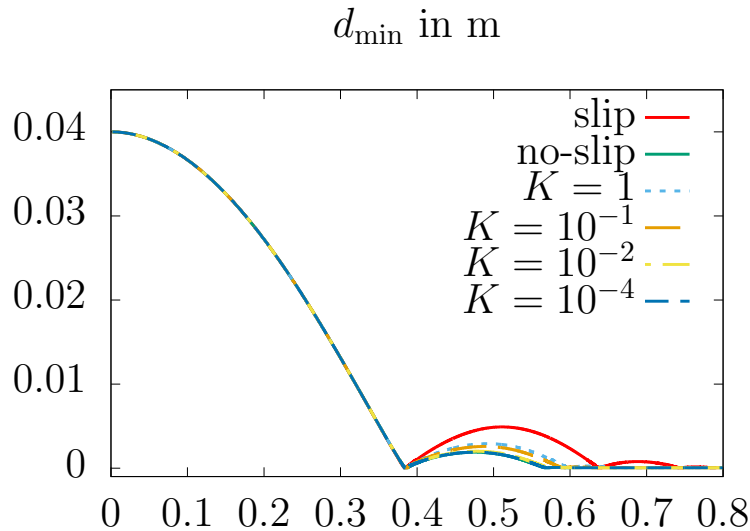
$$\delta t \in \left[\frac{1}{16\,000}, \frac{1}{500} \right] \text{s}$$

(movie)



Variation of the porous conductivity K

Minimal distance d_{\min} to the ground over time for different conductivities K compared to a relaxation approach (without porous layer) with slip resp. no-slip conditions

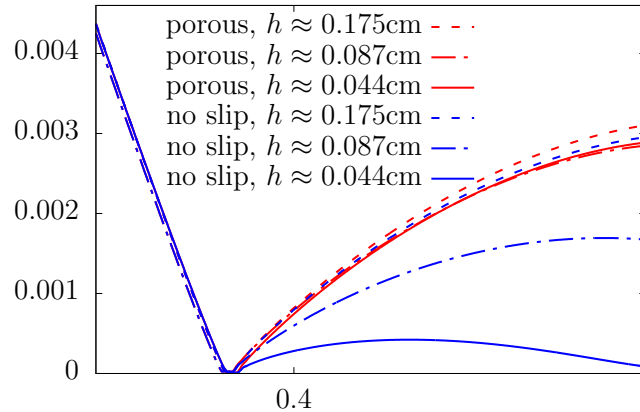


Results with porous layer lie between slip (larger bounce) and no-slip conditions

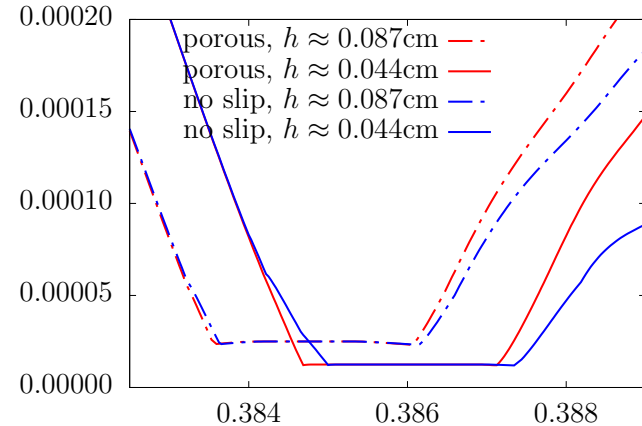
Comparison to relaxation without porous layer

Minimal distance d_{\min} and zoom-in for different h

d_{\min} in m (zoom 1)



d_{\min} in m (zoom 2)



Solid velocities \dot{d}_y
at impact (t_i) and
release (t_r)

h	With porous layer		Only no-slip	
	$-\overline{\dot{d}_y}(t_i)$	$\overline{\dot{d}_y}(t_r)$	$-\overline{\dot{d}_y}(t_i)$	$\overline{\dot{d}_y}(t_r)$
$1.75 \cdot 10^{-3}$	$1.12 \cdot 10^{-1}$	$8.87 \cdot 10^{-2}$	$1.11 \cdot 10^{-1}$	$8.81 \cdot 10^{-2}$
$8.77 \cdot 10^{-2}$	$1.05 \cdot 10^{-1}$	$8.48 \cdot 10^{-2}$	$9.74 \cdot 10^{-2}$	$7.59 \cdot 10^{-2}$
$4.39 \cdot 10^{-2}$	$1.03 \cdot 10^{-1}$	$8.91 \cdot 10^{-2}$	$7.25 \cdot 10^{-2}$	$5.98 \cdot 10^{-2}$

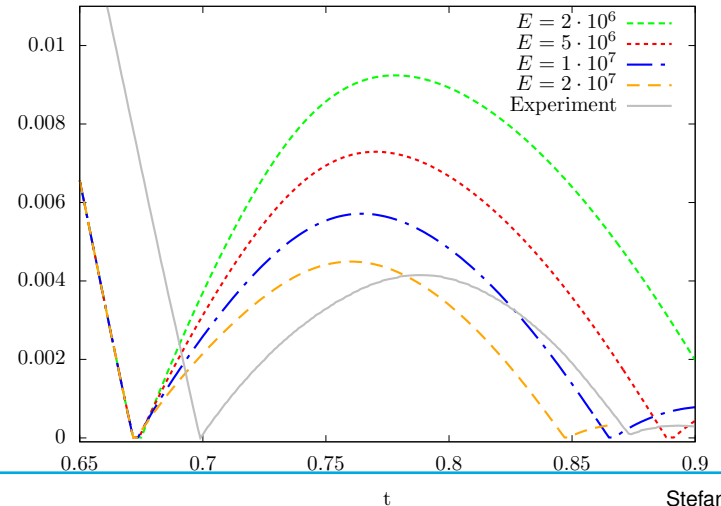
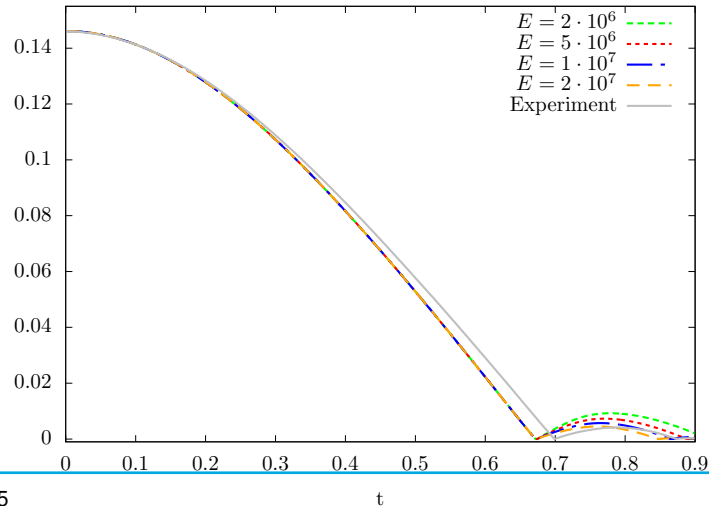
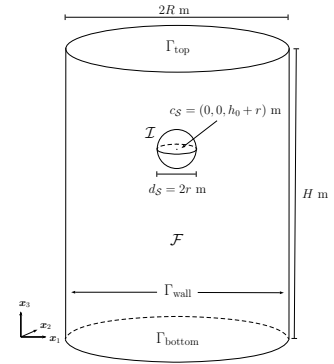
Comparison with experimental benchmark (2.5d)

Falling rubber ball within water-glycerine mixture

(HAGEMEIER, THEVENIN, RICHTER, Int J Multiphase Flow (2020),

VON WAHL, RICHTER, F., HAGEMEIER, Phys Fluids (2021))

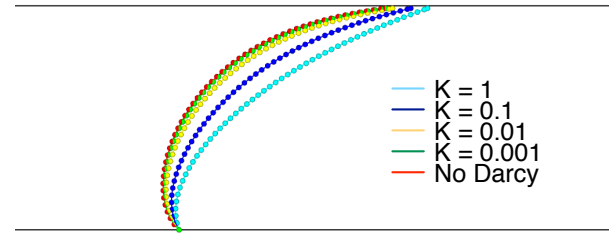
- Solid parameters $\nu = 0.4999$, $E \in [1.7 \cdot 10^6, 2.1 \cdot 10^7]$
- Minimal distance d_{\min} and zoom-in for different Young's moduli E :



Conclusion and outlook

Conclusion

- Mechanically consistent and easy implementable model for FSI and contact considering seepage
- Applied also to a thin elastic solid (beam model) in a mixed coordinate framework with unfitted finite elements (BURMAN, FERNÁNDEZ, F., GEROSA, CMAME 2021)



Outlook

- Comparison of different numerical approaches for benchmark configuration (2d/2.5d)

Main references

- E. Burman, M.A. Fernández, S. Frei: A Nitsche-based formulation for fluid-structure interactions with contact, ESAIM M2AN 54(2), 531-564 (2020)
- S. Frei, F.M. Gerosa, E. Burman, M.A. Fernández: A mechanically consistent model for fluid-structure interactions with contact including seepage, Comp Methods Appl Mech Eng 392, 114637 (2022)
- H. von Wahl, T. Richter, S. Frei, T. Hagemeyer: Falling balls in a viscous fluid with contact: Comparing numerical simulations with experimental data, Phys Fluids 33, 033304 (2021)