

On different constitutive relations and boundary conditions for fluids

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Outline

- implicitly constituted incompressible fluid problem
- typical problems, motivation
- our result on existence and uniqueness
- numerical analysis - ongoing work

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Bulíček, Málek, M: *On nonlinear parabolic problems with implicit constitutive equations involving flux*, M3AS (2021).

Bulíček, Málek, M: *On unsteady internal flows of incompressible fluids characterized by implicit constitutive equations in the bulk and on the boundary*, J. Math. Fluid Mech. (2023).

Gazca Orozco, Gmeineder, MK, Tscherpel, *in preparation 2024*.

Navier–Stokes-like problem

- flow of homogeneous, incompressible fluids

$$\begin{aligned}\partial_t \mathbf{v} + \operatorname{div}(\mathbf{v} \otimes \mathbf{v}) - \operatorname{div} \mathbf{S} + \nabla p &= \mathbf{f} && \text{in } Q := (0, T) \times \Omega, \\ \operatorname{div} \mathbf{v} &= 0 && \text{in } Q, \\ \mathbf{v} \cdot \mathbf{n} &= 0 && \text{on } \Gamma := (0, T) \times \partial\Omega, \\ \mathbf{v}(0) &= \mathbf{v}_0 && \text{in } \Omega,\end{aligned}\tag{1}$$

- Navier–Stokes: $\mathbf{S} = 2\nu D\mathbf{v}$ in $Q \longrightarrow$ implicit: $\mathcal{G}(\mathbf{S}, D\mathbf{v}) = 0$ in Q
- boundary condition: $\mathbf{s} := -(\mathbf{S}\mathbf{n})_\tau = \gamma \mathbf{v}_\tau$ on $\Gamma \longrightarrow$ implicit: $\mathbf{g}(\mathbf{s}, \mathbf{v}_\tau) = 0$ on Γ

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- for example

$$\mathbf{S} - |\mathbf{D}\mathbf{v}|^{r-2} \mathbf{D}\mathbf{v} = 0 \quad \text{or} \quad \mathbf{D}\mathbf{v} - (|\mathbf{S}| - \sigma_*)^+ \frac{\mathbf{S}}{|\mathbf{S}|} = 0$$

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- weak formulation:

$$\langle \partial_t \mathbf{v}, \boldsymbol{\varphi} \rangle - \int_{\Omega} (\mathbf{v} \otimes \mathbf{v}) : \nabla \boldsymbol{\varphi} \, d\mathbf{x} + \int_{\Omega} \mathbf{S} : \mathbf{D}\boldsymbol{\varphi} \, d\mathbf{x} + \int_{\partial\Omega} \mathbf{s} \cdot \boldsymbol{\varphi} \, dS = \langle \mathbf{f}, \boldsymbol{\varphi} \rangle,$$

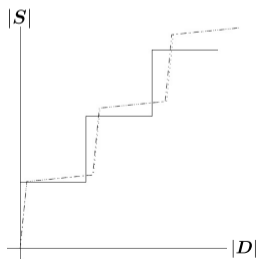
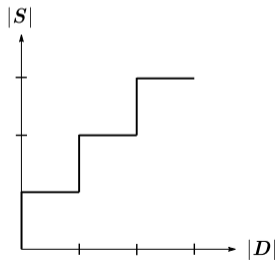
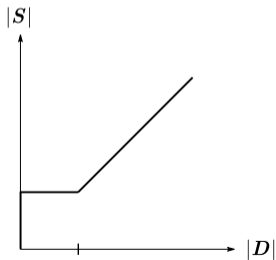
Power-law and/or activating effect

$r \rightarrow \infty$	rigid body		limiting shear-rate		Euler/limiting shear-rate	
$r \in (2, \infty)$	rigid/shear-thickening		shear-thickening		Euler/shear-thickening	
$r = 2$	Bingham = rigid/Navier-Stokes		Navier-Stokes	$ S $ 	Euler/Navier-Stokes	
$r \in (1, 2)$	rigid/shear-thinning		shear-thinning		Euler/shear-thinning	
$r \rightarrow 1$	perfect plastic		limiting shear stress		Euler	
	$ S \leq \sigma_* \iff Dv = O$		no activation		$ Dv \leq \delta_* \iff S = O$	

Blechta, Málek, Rajagopal ('20). Power-law $S = |Dv|^{r-2} Dv$; activation $Dv = (|S| - \sigma_*)^+ \frac{S}{|S|}$.

Implicit constitutive theory

- Minty (1962): monotone mappings theory
- Rajagopal (2003, 2006): physical motivation and characterization of implicit theory
- Bulíček, Gwiazda, Málek, Świerczewska-Gwiazda (2009, 2012): mathematical analysis for graphs with selection: $(S, Dv) \in \mathcal{A}$
proof: approximation by convolution
- Bulíček, Málek, M. (2021, 2023): $\mathcal{G}(S, Dv) = 0$, no selection needed, algebraical approximation, easy-to-verify assumptions



Maximal monotone r -graph $\mathcal{A} \subset \mathbb{R}^{d \times d} \times \mathbb{R}^{d \times d}$, $r \in (1, \infty)$, $r' := \frac{r}{r-1}$

(A1) $(0, 0) \in \mathcal{A}$,

(A2) **monotonicity**: for any $(\mathbf{S}_1, \mathbf{D}_1), (\mathbf{S}_2, \mathbf{D}_2) \in \mathcal{A}$,

$$(\mathbf{S}_1 - \mathbf{S}_2) : (\mathbf{D}_1 - \mathbf{D}_2) \geq 0,$$

(A3) **maximality**: if for some (\mathbf{S}, \mathbf{D}) and all $(\bar{\mathbf{S}}, \bar{\mathbf{D}}) \in \mathcal{A}$

$$(\mathbf{S} - \bar{\mathbf{S}}) : (\mathbf{D} - \bar{\mathbf{D}}) \geq 0$$

holds, then $(\mathbf{S}, \mathbf{D}) \in \mathcal{A}$,

(A4) **r -growth & coercivity**: there exist $C_1, C_2 > 0$ such that for all $(\mathbf{S}, \mathbf{D}) \in \mathcal{A}$ there holds

$$\mathbf{S} : \mathbf{D} \geq C_1(|\mathbf{S}|^{r'} + |\mathbf{D}|^r) - C_2.$$

Implicit function $\mathcal{G} : \mathbb{R}^{d \times d} \times \mathbb{R}^{d \times d} \rightarrow \mathbb{R}^{d \times d}$

(G1) $\mathcal{G}(\mathbf{S}, \mathbf{D}) \in \mathcal{C}^{0,1}$ and $\mathcal{G}(0, 0) = 0$,

(G2) for almost all (\mathbf{S}, \mathbf{D}) :

$$\mathcal{G}_{\mathbf{S}} \geq 0, \quad \mathcal{G}_{\mathbf{D}} \leq 0, \quad \mathcal{G}_{\mathbf{D}}(\mathcal{G}_{\mathbf{S}})^T \leq 0, \quad \mathcal{G}_{\mathbf{S}} - \mathcal{G}_{\mathbf{D}} > 0,$$

(G3)

either $\forall \mathbf{D} \quad \liminf_{|\mathbf{S}| \rightarrow +\infty} \mathcal{G}(\mathbf{S}, \mathbf{D}) : \mathbf{S} > 0$

or $\forall \mathbf{S} \quad \limsup_{|\mathbf{D}| \rightarrow +\infty} \mathcal{G}(\mathbf{S}, \mathbf{D}) : \mathbf{D} < 0,$

(G4) there exist $c_1, c_2 > 0$ such that for all $\mathcal{G}(\mathbf{S}, \mathbf{D}) = 0$ we have

$$\mathbf{S} : \mathbf{D} \geq c_1(|\mathbf{S}|^{r'} + |\mathbf{D}|^r) - c_2.$$

Existence of weak solution

Denote $z := \max\{r, q, \frac{(d+2)r}{(d+2)r-2d}\}$. For any

- $\Omega \in \mathcal{C}^{0,1}$, $T > 0$, $\mathbf{v}_0 \in H$, $\mathbf{f} \in L^{r'}(0, T; (W_{\mathbf{n}, \text{div}}^{1,r}(\Omega))^*)$
- \mathcal{G} satisfying (G1) – (G4) with $r > \frac{2d}{d+2}$,
- \mathbf{g} satisfying (g1) – (g4) with $q > 1$,

there exist $(\mathbf{v}, \mathbf{S}, \mathbf{s})$ a weak solution to the problem (1) with implicit relations

$$\begin{aligned} \mathcal{G}(\mathbf{S}, D\mathbf{v}) &= 0 && \text{in } Q, \\ \mathbf{g}(\mathbf{s}, \mathbf{v}_\tau) &= 0 && \text{on } \Gamma, \end{aligned}$$

such that

$$\begin{aligned} \mathbf{v} &\in L^r(0, T; W_{\mathbf{n}, \text{div}}^{1,r}(\Omega)) \cap \mathcal{C}_w([0, T]; L_{\mathbf{n}, \text{div}}^2(\Omega)) \cap L^q(\Gamma), \\ \partial_t \mathbf{v} &\in L^{z'}(0, T; (W_{\mathbf{n}, \text{div}}^{1,z}(\Omega))^*), \\ \mathbf{S} &\in L^{r'}(Q), \quad \mathbf{s} \in L^{q'}(\Gamma). \end{aligned}$$

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Also, if $\Omega \in \mathcal{C}^{1,1}$ and $\mathbf{f} \in L^{r'}(0, T; (W_{\mathbf{n}}^{1,r}(\Omega))^*)$, there exists a pressure $p \in L^{z'}(0, T; L^{z'}(\Omega))$.

On uniqueness

+ parabolic, no convective term: $r > 1$ (Bulíček, Málek, Maringová 2021)

+ power-law-fluid:

$$r \geq \frac{d+2}{2} \quad (\text{Ladyzhenskaya 1960s})$$

$$r \geq \frac{11}{5}, \quad d = 3 \quad \& \quad \text{add. growth cond.} \quad (\text{Bulíček, Kaplický, Pražák 2019})$$

+ implicit fluids:

$$r \geq \frac{3d+2}{d+2} \quad (\text{Bulíček, Málek, Maringová 2023})$$

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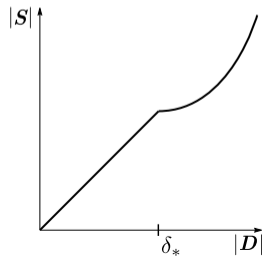
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Navier-Stokes prior to activation,

i.e. Navier-Stokes for $|\mathbf{D}\mathbf{v}| \leq \delta_*$, $\delta_* > 0$ arbitrary
and power-law for $|\mathbf{D}\mathbf{v}| > \delta_*$



On non-uniqueness

- for very weak solutions

$$r < \frac{3d+2}{d+2} \quad (\text{Burczak, Modena, Székelyhidi 2021})$$

- for Leray solutions $r = 2$ (Albritton, Brué, Colombo 2022)

Numerical analysis - boundary conditions

BMM '21: parabolic problem, $\mathcal{G}(\mathbf{S}, D\mathbf{v})$, $r > 1$;

BMM '23: Navier–Stokes-like, $\mathcal{G}(\mathbf{S}, D\mathbf{v})$, $r > \frac{2d}{d+2}$; $\mathbf{g}(\mathbf{s}, \mathbf{v})$, $q > 1$;

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GO-G-MK-T \sim '24: Navier–Stokes, $\mathbf{S} = 2\nu\mathbf{D}\mathbf{v}$, $r = 2$; different BC, $\mathbf{g}(\mathbf{s}, \mathbf{v})$, $1 \leq q < \infty$.

- existence of discrete and weak solution
- stationary / time dependent problem
- explicit / implicit BC
- coercive / non-coercive BC
- monotone / non-monotone BC

Numerical analysis - boundary conditions

In particular:

- Nitsche penalisation due to approximation of the domain - cannot impose the BC directly
- Korn inequality with normal traces (under geometric assumption on the domain)

$$\|\nabla \mathbf{v}\|_{L^p(\Omega)} \lesssim \|\mathbf{D}\mathbf{v}\|_{L^p(\Omega)} + \|\mathbf{v} \cdot \mathbf{n}\|_{L^p(\Gamma)} \quad \text{where } \Gamma \subset \partial\Omega$$

(see Bauer, Pauly 2016 for $\mathbf{v} \cdot \mathbf{n} = 0$)

- Tresca slip: $q = 1$



- (explicit) non-monotone BC:

$$(\mathbf{s}^*(\mathbf{v}_1) - \mathbf{s}^*(\mathbf{v}_2)) \cdot (\mathbf{v}_1 - \mathbf{v}_2) \geq -\lambda |\mathbf{v}_1 - \mathbf{v}_2|^2 \quad \text{for some } \lambda \geq 0$$

- dynamic slip (see Abbatiello, Bulíček, M. 2021):

$$\mathbf{s} = \alpha \mathbf{v} + \beta \partial_t \mathbf{v} \quad \text{for some } \alpha, \beta \geq 0$$