

Towards a sharper phase-field method for microstructure evolution problems

Stanisław Stupkiewicz

joint work with Jędrzej Dobrzański

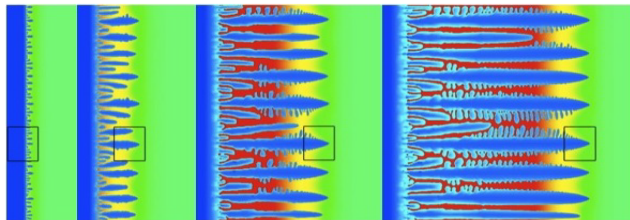
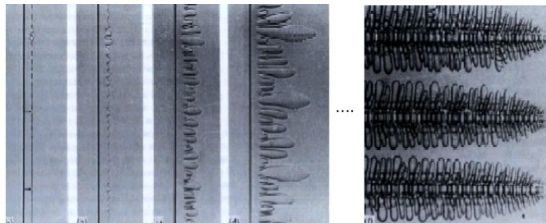


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Modelling, PDE analysis and computational mathematics in materials science
MPDE, Prague, Czech Republic, 22-27/09/2024

Microstructure evolution

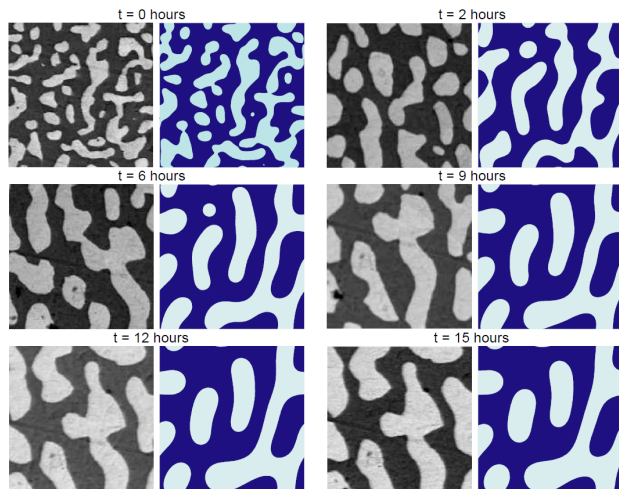
- Directional solidification, dendritic growth, ...



[Provatas & Elder, 2010]

Microstructure evolution

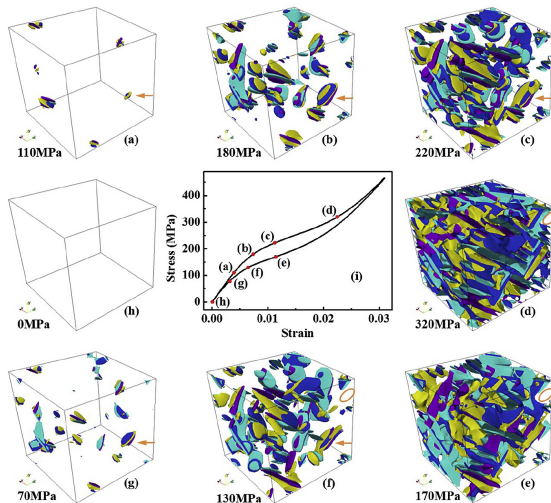
- Spinodal decomposition – microstructure coarsening by diffusion



Thermal ageing of tin–lead solder at 150° [Ubachs, Schreurs & Geers, *JMPS* 2004]

Microstructure evolution

- Martensitic transformation



[Zhu et al, 2017]

Phase-field method: diffuse-interface modelling approach

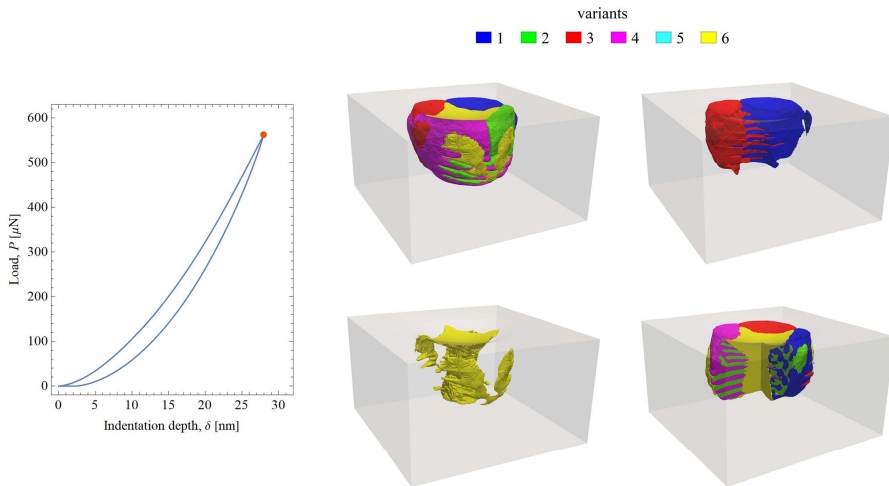
- tracking of interfaces is avoided (interfaces are diffuse)
- microstructure evolution is simulated on a fixed (**but fine!**) computational grid
- various physical phenomena can be included

Some literature references on the phase-field method:

- ▶ general reviews [Chen, 2002; Moelans et al, 2008; Steinbach, 2009; Provatas & Elder, 2010; Wang & Li, 2010]
- ▶ phase-field method for martensitic transformation [Wang & Khachaturyan, 1997; Artemev et al, 2000; Wen et al, 2000; Levitas & Preston, 2002; Ahluwalia et al, 2008; Shu & Yen, 2008; Lei et al, 2010; Hildebrand & Miehe, 2012; Borukhovich et al, 2014; Schoof et al, 2019; Liu et al, 2024, ...]
- ▶ finite-element-based finite-strain phase-field models [Levitas et al, 2009; Clayton & Knap, 2011; Hildebrand & Miehe, 2012; Mosler et al, 2014; Tîma et al, 2016, 2018, 2021; Bartels & Mosler, 2017; Kiefer et al, 2017; Basak & Levitas, 2019; Rezaee-Hajidehi & Stupkiewicz, 2020, 2021]

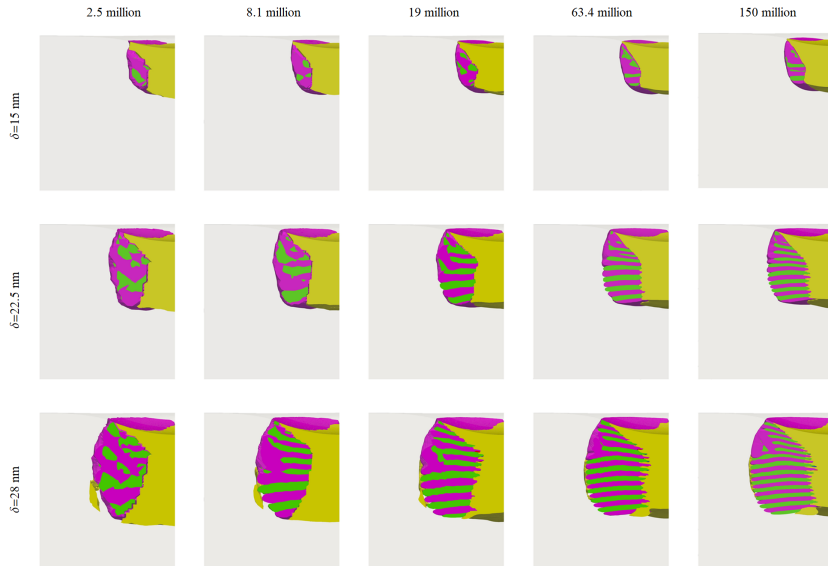
Martensitic transformation in CuAlNi during nano-indentation

[Tůma, Rezaee-Hajidehi, Hron, Farrell & S, *CMAME* 2021]



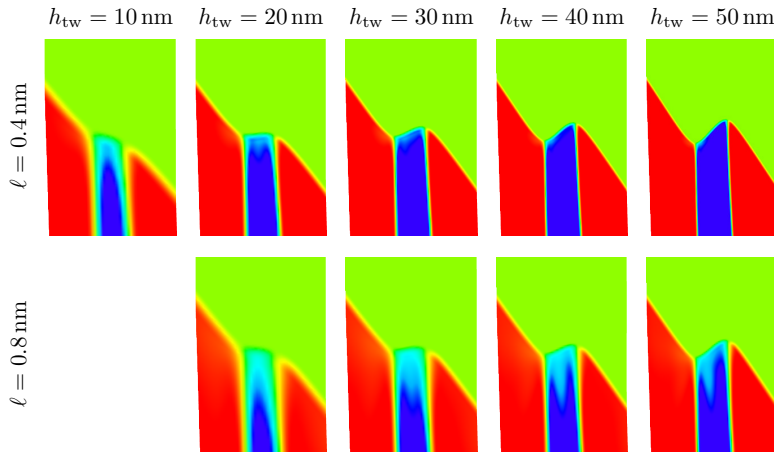
Fine mesh is needed $\rightarrow$ large-scale problems

$\ell = \text{const}$



[Tůma et al, 2021]

So why not to increase ℓ ?



[Tůma & S, *IJSS* 2016]

Aim of our work

- Can we improve computational efficiency of the phase-field method?
 - ▶ sharper interfaces $\rightarrow$ larger element size $\rightarrow$ reduced computational cost
 - ▶ keeping simplicity (e.g., avoid XFEM)

$\rightarrow$ sharp phase-field method [Finel et al, 2018] – but relying on the finite difference method

Outline

- 1 Conventional phase field method
- 2 LET: lamination-based treatment of weak discontinuities
- 3 LET-PF: hybrid diffuse-semisharp approach for evolving interfaces

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Prototype phase-field model

- Assumptions:
 - two phases
 - viscous evolution
- Constitutive framework for each phase: elasticity with eigenstrain ($i = 1, 2$)

$$F_i(\boldsymbol{\varepsilon}) = F_i^0 + \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_i^t) \cdot \mathbb{L}_i(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_i^t), \quad \boldsymbol{\sigma} = \mathbb{L}_i(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_i^t)$$

- Phase transformation: continuous order parameter $\phi \approx$ volume fraction ($0 \leq \phi \leq 1$)
- Bulk free energy density

$$F_B(\boldsymbol{\varepsilon}, \phi) = F^0 + \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^t) \cdot \mathbb{L}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^t)$$

Interpolated properties:

$$\begin{aligned}\boldsymbol{\varepsilon}^t(\phi) &= (1 - \phi)\boldsymbol{\varepsilon}_1^t + \phi\boldsymbol{\varepsilon}_2^t \\ \mathbb{L}(\phi) &= (1 - \phi)\mathbb{L}_1 + \phi\mathbb{L}_2 \\ F^0(\phi) &= (1 - \phi)F_1^0 + \phi F_2^0\end{aligned}$$

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Interpolated properties:

$$\begin{aligned} \boldsymbol{\varepsilon}^t(\phi) &= (1 - h(\phi))\boldsymbol{\varepsilon}_1^t + h(\phi)\boldsymbol{\varepsilon}_2^t & h(0) = 0, \quad h(1) = 1, \quad h'(0) = h'(1) = 0 \\ \mathbb{L}(\phi) &= (1 - h(\phi))\mathbb{L}_1 + h(\phi)\mathbb{L}_2 \\ F^0(\phi) &= (1 - h(\phi))F_1^0 + h(\phi)F_2^0 \end{aligned}$$

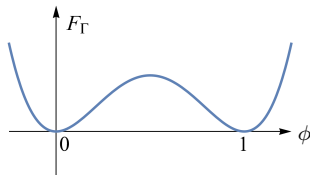
Prototype phase-field model (2)

- Bulk free energy density

$$F_B(\boldsymbol{\varepsilon}, \phi) = F^0(\phi) + \frac{1}{2}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^t(\phi)) \cdot \mathbb{L}(\phi)(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^t(\phi))$$

- Interfacial free energy density (so-called double-well potential)

$$F_\Gamma(\phi, \nabla\phi) = \gamma \left(\frac{6}{\ell} \phi^2(1-\phi)^2 + \frac{3\ell}{2} |\nabla\phi|^2 \right)$$



- Total free energy density

$$F(\boldsymbol{\varepsilon}, \phi, \nabla\phi) = F_B(\boldsymbol{\varepsilon}, \phi) + F_\Gamma(\phi, \nabla\phi)$$

- Dissipation potential (of viscous type)

$$D(\dot{\phi}) = \frac{1}{2m} \dot{\phi}^2$$

+ variational framework $\rightarrow$ complete model

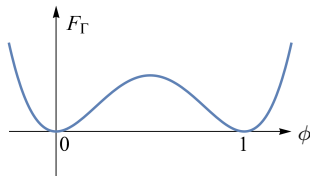
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Variational framework

- Potential energy functional

$$\mathcal{E}[\mathbf{u}, \phi] = \mathcal{F}[\mathbf{u}, \phi] + \Omega[\mathbf{u}], \quad \mathcal{F}[\mathbf{u}, \phi] = \int_B F(\nabla_s \mathbf{u}, \phi, \nabla \phi) \, dV$$

→ $\Omega[\mathbf{u}]$ – potential of external loads, e.g., $\Omega[\mathbf{u}] = - \int_{\partial_t B} \mathbf{t}^* \cdot \mathbf{u} \, dS$

- Global rate potential

$$\Pi[\dot{\mathbf{u}}, \dot{\phi}; \mathbf{u}, \phi] = \dot{\mathcal{E}}[\dot{\mathbf{u}}, \dot{\phi}; \mathbf{u}, \phi] + \mathcal{D}[\dot{\phi}], \quad \mathcal{D}[\dot{\phi}] = \int_B D(\dot{\phi}) \, dV$$

- Minimization problem [Miehe, 2011; Hildebrand & Miehe, 2012; Tuma, Petryk & S, 2016, 2018]

$$\{\dot{\mathbf{u}}, \dot{\phi}\} = \arg \min_{\dot{\mathbf{u}}, \dot{\phi}} \Pi[\dot{\mathbf{u}}, \dot{\phi}; \mathbf{u}, \phi]$$

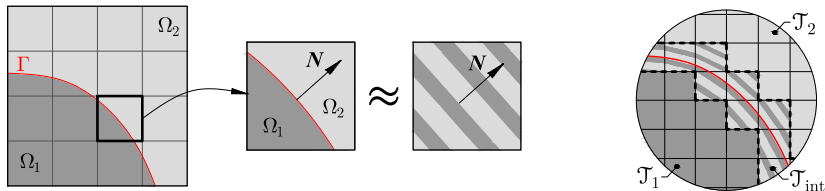
- ▶ $\delta_{\dot{\mathbf{u}}} \Pi = 0 \rightarrow$ equilibrium equation (virtual work principle)
- ▶ $\delta_{\dot{\phi}} \Pi = 0 \rightarrow$ evolution equation for ϕ

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Laminated element technique (LET)

[Dobrzański, Wojtacki & S, *Comput. Struct.* 2024]



- Interfaces defined implicitly by a level set function ϕ
- Total free energy (FE discretization + LET)

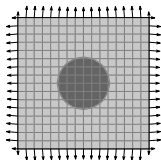
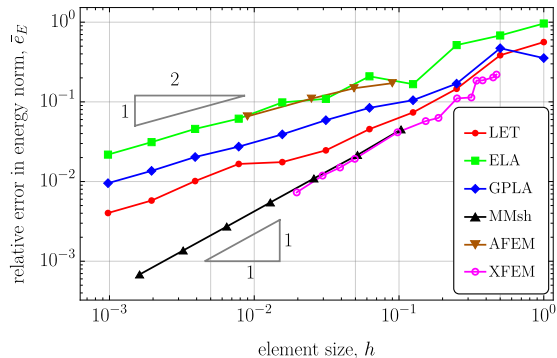
$$\mathcal{F} = \mathcal{F}_B = \sum_{i=1}^2 \left(\sum_{\omega \in \mathcal{T}_i} \int_{\omega} F_i \, dV \right) + \sum_{\omega \in \mathcal{T}_{\text{int}}} \int_{\omega} \overline{F} \, dV$$

- ▶ microstructure of the laminate $(\eta, \mathbf{N})$ depends on the position of the interface within the element
- ▶ overall (homogenized) bulk energy of simple laminates within laminated elements, $\omega \in \mathcal{T}_{\text{int}}$

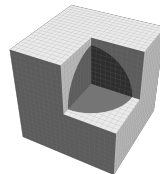
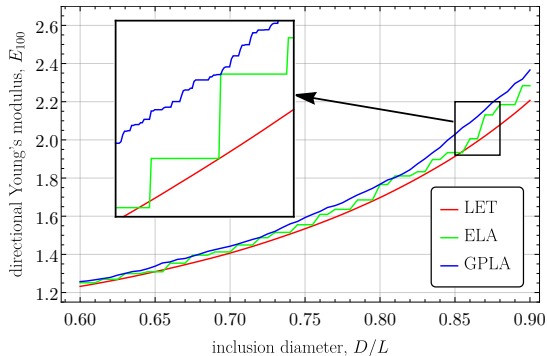
$$\overline{F} = \overline{F}(\varepsilon, \eta, \mathbf{N}) = (1 - \eta)F_1 + \eta F_2$$

LET performance: elastic inclusion problem

convergence with mesh refinement



response as a function of radius



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LET-PF: combining the phase-field method with LET

- Total free energy

$$\mathcal{F} = \mathcal{F}_B + \mathcal{F}_\Gamma$$

- **Interfacial** free energy $\rightarrow$ **diffuse** treatment of interfaces by the **phase-field (PF)** method

$$\mathcal{F}_\Gamma = \sum_{\omega \in \mathcal{T}} \int_{\omega} F_\Gamma \, dV, \quad F_\Gamma = \gamma \left(\frac{6}{\ell} \phi^2 (1 - \phi)^2 + \frac{3\ell}{2} |\nabla \phi|^2 \right)$$

- **Bulk** free energy $\rightarrow$ **semisharp** treatment of interfaces by **LET**

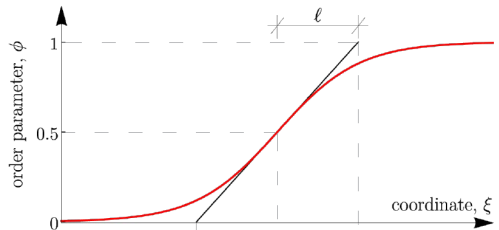
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- ▶ the order parameter ϕ plays the role of a level set function ($\phi = \frac{1}{2} \rightarrow$ interface)

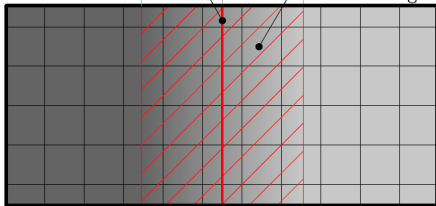
- Variational framework

$$\Pi = \dot{\mathcal{F}} + \mathcal{D} \rightarrow \min$$

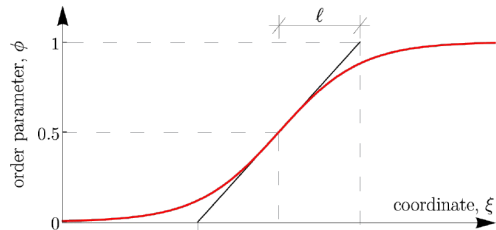
LET-PF: combining the phase-field method with LET



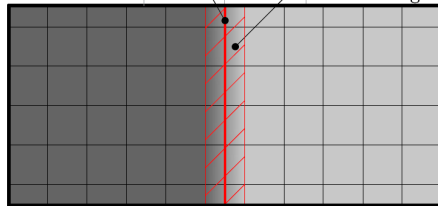
real interface $h \leq \ell/2$ bulk energy transition region



conventional phase-field



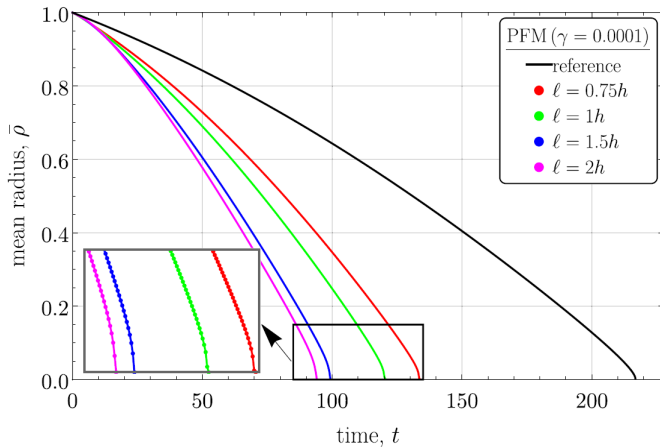
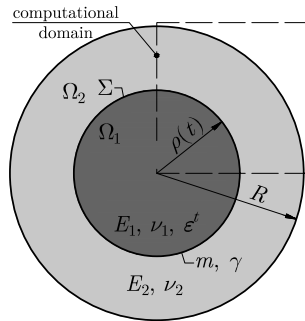
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LET-PF

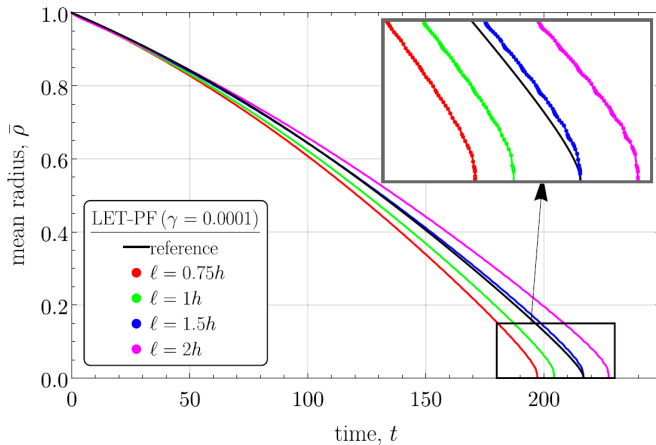
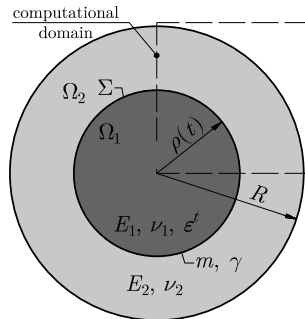
Evolving (vanishing) circular inclusion: conventional phase-field method

- evolution driven by elastic strain energy, large time increments, coarse mesh (2D)
- analytical solution of the reference 1D sharp-interface problem



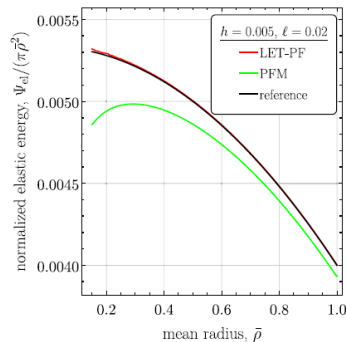
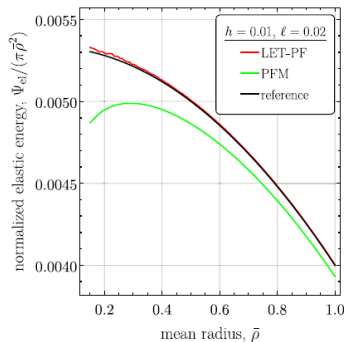
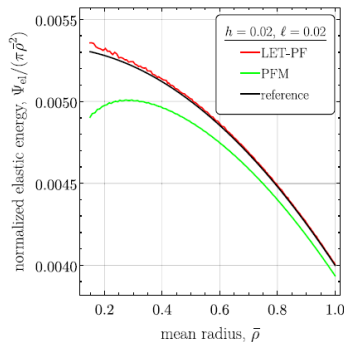
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Evolving (vanishing) circular inclusion: effect of mesh refinement

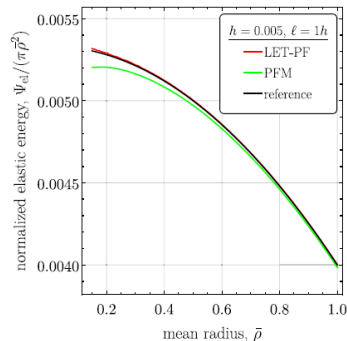
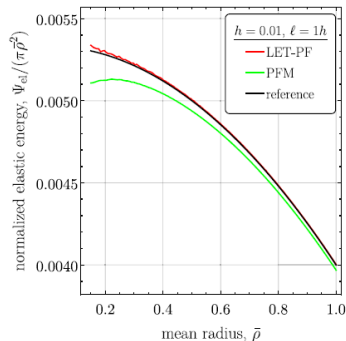
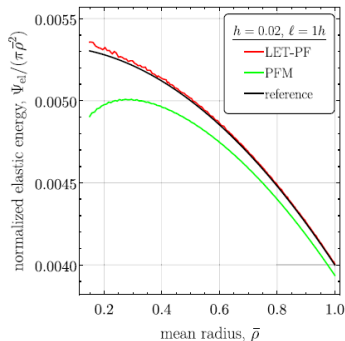
- elastic strain energy as a function of inclusion radius, $\ell = \text{const}$



- error due to the finite thickness of the diffuse interface in the conventional phase-field method
- the error is reduced thanks to the semisharp LET method

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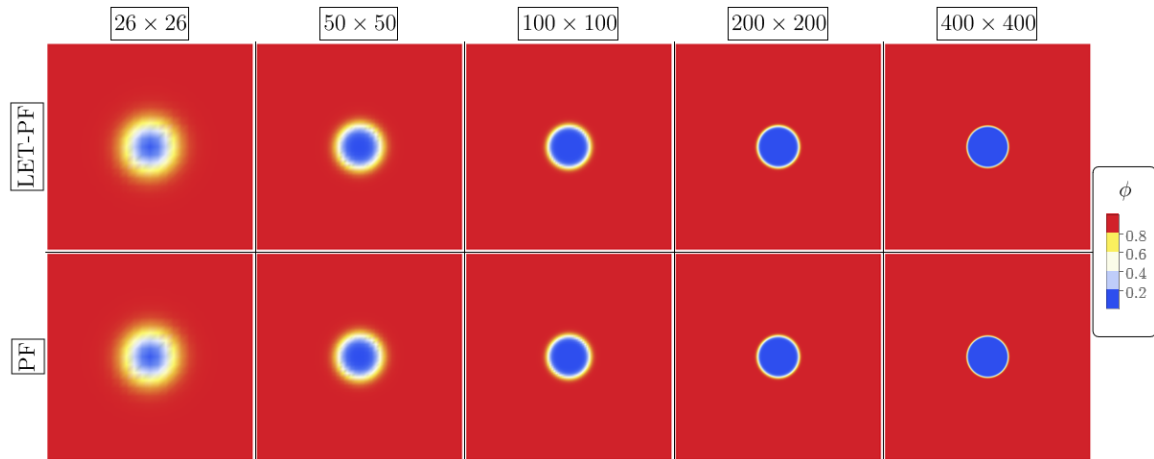
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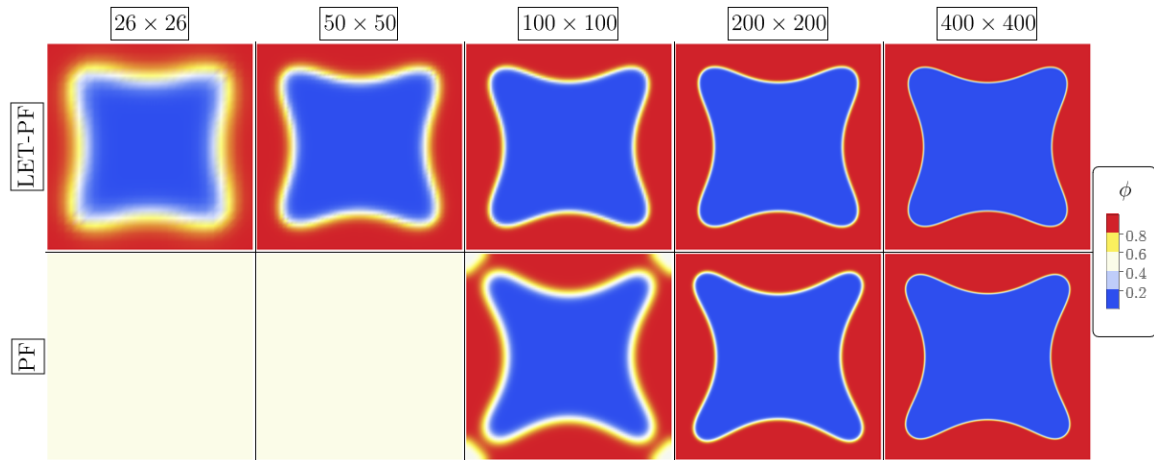
Evolving inclusion in a constrained domain: effect of mesh size

- LET-PF vs. conventional phase-field (PF), $h/\ell = \text{const}$ – initial state



Evolving inclusion in a constrained domain: effect of mesh size

- LET-PF vs. conventional phase-field (PF), $h/\ell = \text{const}$ – final state



Evolving inclusions in a constrained domain: effect of mesh size

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Summary

In short:

- *Proof of concept* of a new approach to phase-field modelling of microstructure evolution

In more detail:

- LET-PF, a hybrid diffuse-semisharp approach combining
 - ① phase-field method – interfacial energy contribution
 - ② laminated element technique (LET) – bulk energy contribution
- Features:
 - simple (element-based implementation)
 - for the problems studied, LET-PF exhibits higher accuracy than the conventional phase-field method

Even more details:

J. Dobrzański, K. Wojtacki & S. Stupkiewicz, Lamination-based efficient treatment of weak discontinuities for non-conforming finite element meshes, *Comput. Struct.* **291**, 107209, 2024.

J. Dobrzański & S. Stupkiewicz, Towards a sharper phase-field method: A hybrid diffuse-semisharp approach for microstructure evolution problems, *Comp. Meth. Appl. Mech. Eng.* **423**, 116841, 2024.

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